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1	Constitutive model for soil-rock mixtures in the light of an
2	updated skeleton void ratio concept
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21 Abstract: As a type of special geological body, soil-rock mixtures (SRMs) are widely found in 22 nature and used in civil engineering. Many structures, such as rockfill dams, highways and tunnels 23 have used SRMs as building materials. Proper modelling of SRMs is of great importance to 24 capture the complex behavior of this heterogeneous material. In this manuscript, a simple 25 constitutive model incorporating the skeleton void ratio concept is proposed for SRMs with 26 varying soil contents (sc). A prominent feature of the model is a unified description of the 27 behavior of SRMs with varying sc such that only model parameters of pure rock and of pure soil 28 are required. After calibration, the model shows a good capacity to predict the stress-strain 29 response of SRMs under a wide range of sc, void ratios, and confining pressures. In particular, it 30 captures well the non-associated behavior of rock-dominated SRMs with different sc. Furthermore, 31 the sc-value is shown to modify the plastic flow direction of the material without influencing its 32 yield surface.

33 Keywords: Soil-rock mixture; Skeleton void ratio; Soil content; Constitutive modelling; Critical
 34 state line; Flow rule direction; Non-associated behavior

#### 35 **1. Introduction**

Soil-rock mixtures (SRMs) are heterogeneous materials composed of high-strength rocks, fine-grained soils and pores [37, 38, 42]. SRMs are widely encountered in geotechnical engineering, such as natural slopes [6], waste rocks and tailings from mining [8], clay-aggregate mixtures in rockfill dams [12] and tunnels [14]. The soil content (*sc*) is one of the most important factors governing the mechanical behavior of SRMs [3, 26, 29, 33, 36]. In recent years, many field tests, laboratory tests and numerical simulations have revealed that *sc* greatly affects the shear

42	strength [4, 16, 23, 37], failure modes [5, 15, 17], stress-dilatancy [4, 33], and critical state
43	parameters of SRMs [25, 33]. Experimental and numerical results both show that at a low sc, the
44	mechanical behavior of SRMs is primarily governed by inter-granular friction between rock grains
45	While at a high sc beyond a threshold value $(sc)_{th}$ , the mechanical behavior of SRM is primarily
46	governed by friction characteristics of soil grains.
47	Although some pioneering works has been done on the constitutive modelling of SRMs [7,
48	17, 26], properly modeling the behavior of this kind of complex heterogeneous materials remains
49	an open challenge [35, 39, 40]. Predicting the mechanical properties SRMs for varying sc remains
50	a widely open issue.
51	1) In most constitutive models, SRMs with different sc are indeed treated as different
52	materials with their own sc-specific parameters [21, 24, 41]. For example, the shape of the critical
53	state line (CSL) of a SRM and its location depend on sc. Therefore, new laboratory tests are
54	required each time sc is updated in engineering projects, which results in a waste of time and
55	money.
56	2) Some models tried to build empirical equations to link model parameters with sc [23, 35].
57	For instance, some empirical equations have been established to fit CSLs of SRMs with sc based
58	on experimental data. This method usually introduces new parameters into the constitutive model,
59	i.e., parameters in empirical equations. But, these empirical equations lack of physical meaning
60	and their application to SRMs with different lithology, grain shape and gradation of rock and soils

61 remains questionable.

62 3) Although many numerical studies have analyzed the microstructure of SRMs, few63 attempts have been made to incorporate micro- or meso-mechanisms into constitutive models. For

64	example, Wang et al (2021) [29] found that for rock-dominated SRMs, sc does not affect the
65	normal direction of yield surface but it changes the flow rule direction. These findings could be
66	reflected in constitutive models.
67	In soil constitutive models, the global void ratio $e$ has been chosen as one of the main state
68	variable. However, it was found to be an imperfect index to characterize the mechanical behavior
69	of mixed soils like SRMs [29, 33]. This is because such a global index is not able to account for
70	the non-active participation of the soil grains in the force transmission structure within a SRM.
71	Alternatively, skeleton void ratio turns out to be a more appropriate index to reflect the density of
72	SRMs. Skeleton void ratio corresponds to the void ratio of grains constituting the stress-bearing
73	skeleton. In recent studies, this index has shown a strong potential to give unified descriptions of
74	the behavior of SRMs with varying sc [29, 35].
75	The objective of this manuscript is to propose a simple method to predict the stress-strain
76	response of SRMs with varying sc incorporating the updated skeleton void ratio concept. It is
77	organized as follows. First, the updated skeleton void ratio for SRMs proposed by Wang et al.
78	(2022) [29] is briefly reviewed. The advantage of using this skeleton void ratio index to
79	characterize critical state lines of SRMs with varying sc is then shown. Next, a constitutive model
80	incorporating the skeleton void ratio concept is proposed and the model is validated against
81	experimental results. Eventually, effect of sc on the main properties of SRMs is studied and the
82	model capabilities are discussed.

# **2. Review of an updated skeleton void ratio for SRMs**

84 Wang et al. (2022) [29] proposed an updated skeleton void ratio index for SRMs. The main
85 advantage of this updated skeleton void ratio is that it can consider the effect of gradations of both

soils and rocks. The skeleton void ratio proposed by Thevanayagam (2007) [27] is a special case

87 of the updated skeleton void ratio in which mutual interaction between rock and soil grains during

88 packing is neglected.

89 According to Wang et al. (2022) [29], the threshold soil content  $(sc)_{th}$  that separates the 90 rock-dominated structure and soil-dominated structure is:

$$(sc)_{th} = \frac{e_r - b}{1 + e_r + e_s - a - b}$$
(1)

91 where  $e_s$  and  $e_r$  are the minimum void ratios of pure soil grains and pure rock grains, 92 respectively. *a* and *b* are gradation-related parameters:

$$a = (1 + e_s) \exp(\frac{-R_d^{0.5}}{C_{ur}C_{us}})$$
(2)

$$b = e_r \exp(\frac{-R_d}{C_{ur} C_{us}^{0.5}}) \tag{3}$$

93 where  $C_{ur} = \frac{(D_{60})_{rock}}{(D_{10})_{rock}}$  is the coefficient of non-uniformity for the rock fraction,  $C_{us} = \frac{(d_{60})_{soil}}{(d_{10})_{soil}}$  is 94 the coefficient of non-uniformity for the soil fraction and  $R_d = \frac{(D_{50})_{rock}}{(d_{50})_{soil}}$  is grain size disparity 95 ratio.

96 For a rock-dominated structure, i.e.,  $sc < (sc)_{th}$ , its skeleton void ratio  $e_{sk}$  is given by:

$$e_{sk} = \frac{e + sc}{1 - sc} - \frac{sc(1 + e_s)}{1 - sc} \exp(\frac{-R_d^{0.5}}{C_{ur}C_{us}})$$
(4)

97 For a soil-dominated structure, i.e.,  $sc > (sc)_{th}$ , its skeleton void ratio  $e_{sk}$  expresses as:

$$e_{sk} = \frac{e}{sc} - \frac{e_r(1 - sc)}{sc} \exp(\frac{-R_d}{C_{ur}C_{us}^{0.5}})$$
(5)

98 All the parameters introduced in Equations (1)-(5) can be obtained from simple laboratory

99 sieving and compaction tests.

Wang et al. (2022) [29] found that a rock-dominated (or soil-dominated) SRM has similar
stress-strain responses with a pure rock specimen (or a pure soil specimen) if global void ratio of

the pure rock specimen (or pure soil specimen) equals to the skeleton void ratio of the SRM. This
 important finding is essential to unify descriptions of SRMs with varying soil contents.

# 3. Characterization of critical state lines of SRMs with updated skeleton void ratio index

106 The critical state is defined as the state at which the soil continues to deform at constant shear 107 stress and constant volume. It has increasingly been used as a fundamental concept to characterize 108 the strength and deformation properties of soils [2, 20, 22]. Li et al. (1998) [10] found that the

109 critical state lines (CSLs) for cohesionless soils are straight lines in the  $e^{-(p/p_a)^{\xi}}$  plane:

$$e_{cr} = e_{\Gamma} - \lambda (p/p_a)^{\xi} \tag{6}$$

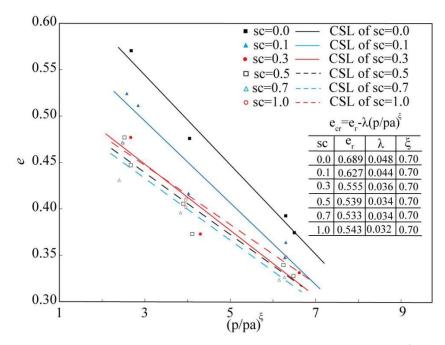
110 Where  $e_{cr}$  is critical state void ratio,  $e_{\Gamma}$  is the theoretical critical void ratio at the atmospheric 111 pressure,  $p_a$  is the atmospheric pressure,  $\lambda$  is the magnitude of the slope, and  $\xi$  is the pressure 112 exponent (with a typical value around 0.7).

Experimental data obtained from conventional drained triaxial tests on SRMs with different soil contents (*sc*=0, 0.1, 0.3, 0.5, 0.7 and 1) under different confining pressures ( $\sigma_3$ =150 kPa, 300 kPa and 600 kPa) conducted by Wang et al. (2022) [29] are adopted here to characterize the CSLs of SRMs with varying *sc* ( $C_{ur}$  = 2.36,  $C_{us}$  = 7.01 and  $R_d$ =5.46). Detailed test procedures can be

117 found in the quoted reference [29].

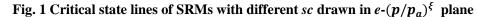
Figure 1 gives the CSLs of SRMs with *sc* ranging from 0 to 1 in *e-p* plane. It can be found that SRMs with different *sc* have different CSLs. The CSL move downward from the pure rock specimen until the soil content reaches  $(sc)_{th}$  (around 0.54), and then the CSL would move upward to the position of pure soil specimen. According to Wang et al. (2022), all the SRMs were prepared with the same global void ratio. For a rock-dominated specimen, the skeleton void ratio

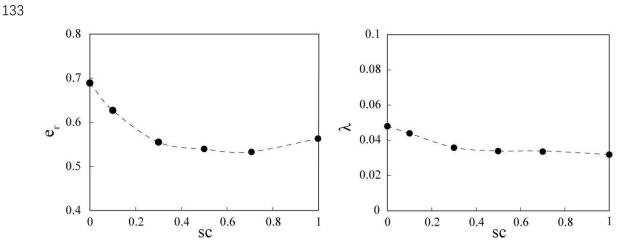
123 increases with *sc* and a more contractant behavior is observed. Therefore, the increase in *sc* widens 124 the gap between current void ratio and critical void ratio and leads to the downward shift of the 125 CSL. On the other hand, for a soil-dominated specimen, the increase in *sc* results in a larger 126 skeleton void ratio and a more dilatant behavior, which consequently narrows the gap. Therefore, 127 CSL will move upwards. Critical state line parameters  $e_{\Gamma}$  and  $\lambda$  are shown in Figure 2.  $e_{\Gamma}$  and 128  $\lambda$  are found to be functions of *sc*, which indicates that SRMs with different *sc* should be treated as 129 different materials. Consequently, each SRM has its own CSL in *e-p* plane.









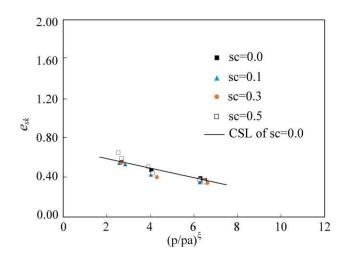


(a)  $e_{\Gamma}$ 

(b) λ

#### Fig. 2 Evolution of critical state parameters with *sc*: (a) $e_{\Gamma}$ ; (b) $\lambda$

134 Critical state data of SRMs with different sc are plotted in  $e_{sk}$ -p plane in Figure 3. Pure rock and pure soil materials have different CSLs because they have different grain shapes and 135 136 gradations. Therefore, for a clearer view, we plot separately the critical state data of rock-dominated structure (Figure 3a) and soil-dominated structure (Figure 3b). It can be seen in 137 Figure 3a, that the CSL of pure rock grains (sc=0) is shared with SRMs of different sc (sc=0, 0.1, 138 139 0.3 and  $0.5 < (sc)_{th} = 0.54)$  for rock-dominated SRMs. Similarly, as shown in Figure 3b, for the 140 soil-dominated specimens (sc=0.7,  $1>(sc)_{th} = 0.54$ ), the CSL of the pure soil (sc=1) is shared with 141 SRM with sc=0.7. Consequently, in the  $e_{sk}$ -p plane, the critical state of all RSM materials is 142 described either by the CSL of the rock or the SCL of the soil depending whether sc is below or above (sc)th. This important finding is essential for establishing the constitutive model, as 143 144 developed in the next section.



(a) Rock-dominated SRMs

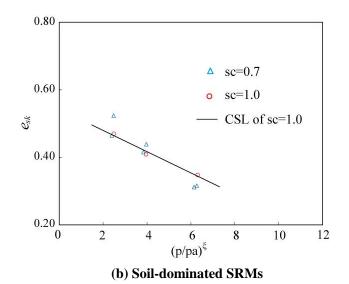


Fig. 1 Critical state lines of SRMs with different *sc* drawn in  $e_{sk} - (p/p_a)^{\xi}$  plane

#### 145 **4.** A simple constitutive model for SRM using the skeleton

#### 146 void ratio index

147 The updated skeleton void ratio concept is incorporated into the state-dependent bounding 148 surface model proposed by Li and Dafalias (2000) [9] to capture the stress-strain behaviors of 149 SRMs with different sc. The fundamental of this model is that the SRMs and its host rock (or soil) 150 with the same skeleton void ratio should exhibit the same stress-strain behavior, which has been 151 reported by many researchers [28, 29, 35]. The fact that rock-dominated (or soil-dominated) 152 SRMs share the same CSL with pure rock (or pure soil) also suggests to build a unified model for 153 SRMs with varying sc from only model parameters of pure rock and pure soil. The constitutive 154 framework is detailed in the following sections.

- 155 **4.1 Elastic behavior**
- 156 According to standard elasto-plasticity, a total strain increment  $d\varepsilon$  is additively split into an
- 157 elastic strain increment  $d\varepsilon^e$  and a plastic strain increment  $d\varepsilon^p$ :

$$\mathbf{d}\boldsymbol{\varepsilon} = \mathbf{d}\boldsymbol{\varepsilon}^{\boldsymbol{e}} + \mathbf{d}\boldsymbol{\varepsilon}^{\boldsymbol{p}} \tag{7}$$

158 The volu

The volumetric and deviatoric elastic strain increments are given respectively by:

$$\mathrm{d}\varepsilon_{\nu}^{e} = \frac{\mathrm{d}p'}{K} \tag{8}$$

$$\mathrm{d}\varepsilon_d^e = \frac{\mathrm{d}q}{3G} \tag{9}$$

159 where *K* is the elastic bulk modulus, *G* is the elastic shear modulus,  $p' = \frac{\sigma_1 + 2\sigma_3}{3}$  is the mean 160 effective stress and  $q = \sigma_1 - \sigma_3$  is the deviatoric stress ( $\sigma_1$ ,  $\sigma_2$ , and  $\sigma_3$  are the principal stress 161 values).

162 *G* can be estimated following [19]:

$$G = G_0 \frac{(2.97 - e)^2}{(1 + e)} \sqrt{p' p_a}$$
(10)

163 where  $G_0$  is an elastic material constant,  $p_a$  is the atmospheric pressure, e is the void of the

164 considered material (soil or rock).

165 The elastic bulk moduli *K* is given by the following relation:

$$K = G \frac{2(1+\nu)}{3(1-2\nu)}$$
(11)

166 where  $\nu$  is the Poison's ratio.

167 The estimation of  $G_0$  and  $\nu$  can be found in Appendix.

#### 168 **4.2 Plastic behavior**

169 For the sake of simplicity, the elasto-plastic model proposed by Li and Dafalias (2000) [9] in

170 the triaxial compression stress space is adopted. In this model, the yield surface is given by:

$$f(p',q,\eta) = q - p'\eta_c = 0$$
 (12)

171 where  $\eta_c$  is the stress ratio when plasticity activates. In the model of Li and Dafalias (2000) [9] 172 the yield surface is assumed to follow the stress state so that  $\eta_c = \eta$ . Note that this hypothesis 173 holds as long as plasticity is activated. In case elastic unloading is considered, the model requires 174 some additional mechanism such as a back-stress [13], or a memory of the reversal stress ratio 175 point [31] which are not considered in the present study.

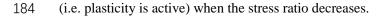
176 According to the theory of plasticity, a loading index dL can be defined as:

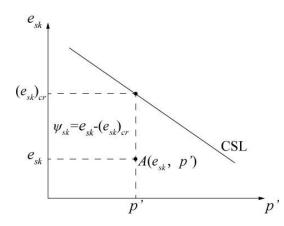
$$dL = \frac{1}{K_p} \left( \frac{\partial f}{\partial p'} dp' + \frac{\partial f}{\partial q} dq \right) = \frac{dq - \eta_c dp'}{K_p} = \frac{p' d\eta}{K_p} \quad \text{since } \eta_c = \eta \tag{13}$$

where  $K_p$  is the plastic hardening modulus and is expressed by Li and Dafalias (2000) [9] as: 177

$$K_p = hG\left(\frac{M}{\eta} - \exp(n\psi)\right) = \frac{hG\exp(n\psi)}{\eta}(M\exp(-n\psi) - \eta)$$
(14)

In the above equation, n is a positive model parameter,  $h = h_1 - h_2 e_0$  with  $h_1$  and  $h_2$  are 178 model parameters and  $e_0$  the initial void ratio, M is the stress ratio at a critical state and 179  $\psi = e - e_{cr}$  is the state parameter defined by Been and Jefferies (1985) [1] with e the current 180 181 void ratio and  $e_{cr}$  the critical state void ratio for the current p'. When using skeleton void ratio instead of the void ratio, a skeleton state parameter  $\psi_{sk}$  is introduced as  $\psi_{sk} = e_{sk} - (e_{cr})_{sk}$ , as 182 183 seen in Figure 4. Note that when softening occurs  $(K_p < 0)$  the plastic multiplier dL is positive





185 186

#### Fig. 4 Definition of skeleton state parameter $\psi_{sk}$

187 Since soil-rock mixtures are typical non-associated materials, the non-associated flow rule is

adopted to define the plastic strain increments, as follows: 188

$$\mathrm{d}\varepsilon_{d}^{p} = \mathrm{d}L = \frac{p'\mathrm{d}\eta}{K_{p}} \tag{15}$$

$$d\varepsilon_{v}^{p} = DdL = \frac{Dp'd\eta}{K_{p}}$$
(16)

189 where  $d\varepsilon_d^p$  and  $d\varepsilon_v^p$  are the plastic deviatoric strain increment and the plastic volumetric strain 190 increment, respectively. The ratio  $D = \frac{d\varepsilon_v^p}{|d\varepsilon_d^p|}$  evaluates the amplitude of the dilatancy and can be

191 chosen as follows:

$$D = \frac{d_0}{M} [M \exp(m\psi) - \eta]$$
(17)

192 where  $d_0$  and *m* are model parameters. For associated flow rule,  $D = \frac{d\varepsilon_v^p}{d\varepsilon_d^p} = \frac{\partial f}{\partial q} = -\eta_c$ .

193 Therefore, for dL>0, the following expressions can be derived for both incremental deviatoric

and volumetric strains:

$$d\varepsilon_d = d\varepsilon_d^e + d\varepsilon_d^p = \frac{dq}{3G} + \frac{P'd\eta}{K_p} = \left(\frac{1}{3G} + \frac{1}{K_p}\right)dq - \frac{\eta}{K_p}dp'$$
(18)

$$d\varepsilon_{\nu} = d\varepsilon_{\nu}^{e} + d\varepsilon_{\nu}^{p} = \frac{dp'}{K} + Dd\varepsilon_{q}^{p} = \frac{D}{K_{p}}dq + \left(\frac{1}{K} - \frac{D\eta}{K_{p}}\right)dp'$$
(19)

Equations (18) and (19) set up the relationship between stain and stress increments.
Eventually, an elastoplastic constitutive relation is derived in a matrix form:

$$\begin{cases} dq \\ dp' \end{cases} = \left[ \begin{pmatrix} 3G & 0 \\ 0 & K \end{pmatrix} - \frac{h(dL)}{K_p + 3G - K\eta D} \begin{pmatrix} 9G^2 & -3KG\eta \\ 3KGD & -K^2\eta D \end{pmatrix} \right] \begin{pmatrix} d\varepsilon_q \\ d\varepsilon_\nu \end{pmatrix}$$
(20)

197 where h(dL) is a Heaviside function with h(L)=1 for dL>0 and h(dL)=0, otherwise (plastic stain 198 exists only when dL is positive). Note that equation (20) is given in the plane (dp', dq) and the 199 dual plane ( $d\varepsilon_v$ ,  $d\varepsilon_q$ ). It can be generalized to any kind of incremental stress and strain tensors by 200 assuming coaxiality between incremental stress and strain.

202 To summarize, the proposed model includes eleven model parameters, all of them being

203 calibrated from drained triaxial tests under different confining pressures: (1) Two elastic

- 204 parameters, i.e.  $G_0$  and  $\nu$ ; (2) Four critical state parameters, i.e. M,  $e_{\Gamma}$ ,  $\lambda$  and  $\xi$ ; (3) Two
- dilatancy parameters, i.e.  $d_0$ , m; (4) Three hardening parameters, i.e.  $h_1$ ,  $h_2$  and n.
- 206 It should be noted that no new parameters have been introduced into this approach. Skeleton

207 void ratio  $e_{sk}$  of SRMs with different *sc* are adopted to replace global void ratio *e* in the 208 above-mentioned equations. Accordingly, skeleton state parameter  $\psi_{sk}$  is used in the place of  $\psi$ . 209 By doing so, the behavior of SRMs with varying *sc* can be predicted from only the model 210 parameters of pure rock and pure soil. Thus, SRMs with different *sc* should no longer be treated as 211 different materials and only require own *sc*-specific model parameters.

#### 212 **5. Model performance**

In this section, the experimental data obtained by Wang et al. (2022) [29] are adopted to test the performance of the constitutive model introduced above. In reference [29], conventional drained triaxial tests on soil-rock mixtures with *sc*=0, 0.1, 0.3, 0.5, 0.7 and 1 were conducted under three different confining pressures ( $\sigma_3$ =150 kPa, 300 kPa and 600 kPa).

217 The threshold soil content  $(sc)_{th}$  of the SRM used in [26] is 0.54. The threshold soil 218 content is larger than that reported by other researchers. The reason may be the use of angular 219 gravel in their tests, which cause large voids among rock grains. Therefore, SRMs with sc=0.1, 0.3220 and 0.5 are rock-dominated SRMs, whose behavior can be predicted from model parameters of 221 pure rock, while stress-strain responses of soil-dominated SRM (i.e., sc=0.7) can be predicted 222 from the model parameters of pure soil. The calibrated model parameters of the pure rock and 223 pure soil used in [26] are listed in Table 1. The description of the calibration process of model 224 parameters is reported in Appendix.

- 225
- 226

Table 1. Model parameters calibrated for pure rock (*sc*=0) and pure soil (*sc*=1).

Elastic	Critical state	Dilatancy	Hardening
 parameters	parameters	parameters	parameters

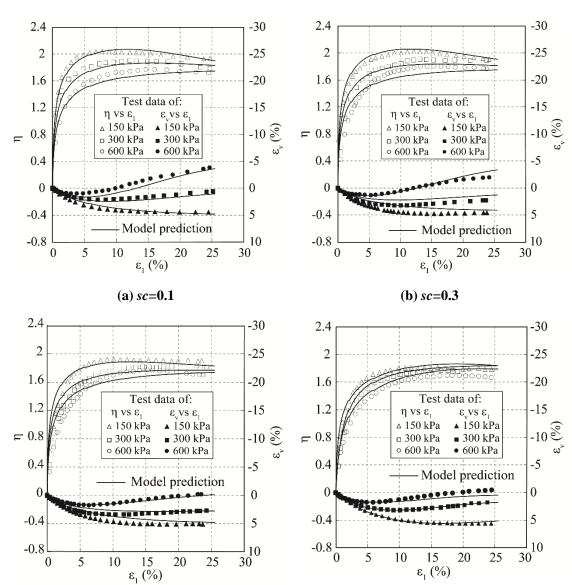
Pure rock ( <i>sc</i> =0)	$G_0 = 200$ $\nu = 0.32$	M = 1.79 $e_{\Gamma} = 0.689$ $\lambda = 0.048$ $\xi = 0.70$	$d_0 = 1.5$ m = -1.0	$h_1 = 0.2$ $h_2 = 0.3$ n = 3.9
Pure soil ( <i>sc</i> =1)	$G_0 = 50$ $\nu = 0.25$	M = 1.70 $e_{\Gamma} = 0.543$ $\lambda = 0.032$ $\xi = 0.70$	$d_0 = 1.8$ m = 0.2	$h_1 = 0.6$ $h_2 = 2.0$ n = 7.0

#### 227 5.1 Stress-strain-volume behavior

Figures 5 and 6 display comparisons between predicted results and experimental results in terms of stress ratios, volumetric strains and evolutions of global void ratios for SRMs with varying *sc*. It can be seen that the stress ratio and the volumetric strain can be fairly described by the model for drained shear responses of SRMs under a range of confining stresses and soil contents.

233 SRMs exhibit post peak strain softening and volumetric expansion at a low confining 234 pressure i.e.,  $\sigma_3=150$  kPa. While at higher confining pressures, i.e.,  $\sigma_3=300$  kPa and 600 kPa, 235 SRMs present strain hardening and the volumetric contraction. As shown in Figure 5, the model 236 can capture both the stress and the volumetric strain behaviors of SRMs with varying sc, e.g., the 237 strain hardening, the volumetric contraction, the strain softening, and the volumetric expansion. In 238 addition, the predicted peak stress ratio, critical stress ratio, phase transformation point from 239 contraction to dilatancy, ultimate volumetric strain, and evolution of void ratio agree well with 240 experimental results.

Some discrepancies between simulated and experimental curves may be observed when *sc* is close to  $(sc)_{th}$ , i.e., when *sc*=0.3 and 0.5. One possible reason for this phenomenon is that when *sc* is close to  $(sc)_{th}$  a dual skeleton structure is formed in SRM which is composed of both rock and



soil grains. Therefore, some discrepancies may occur if we still regard the SRMs as pure

245 rock-skeleton structure or soil-skeleton structure.

244



(d) *sc*=0.7

Fig. 5 Comparison between predicted results and experimental results in terms of stress ratio versus axial strain and volumetric strain versus axial strain.

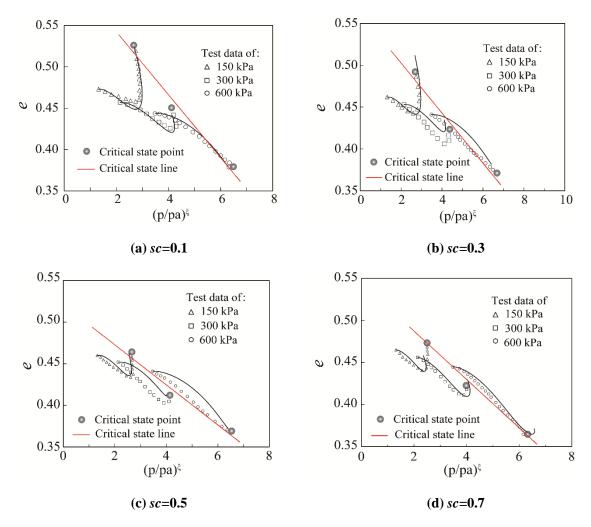


Fig. 6 Comparison between predicted results and experimental results in terms of void ratio versus mean pressure

#### 246 **5.2 Mobilized friction angle**

247 In this section, we assess whether our model can properly reflect the mobilized friction angle

of SRMs with different sc and confining pressures. The mobilized friction angle  $\varphi_m$  is defined as:

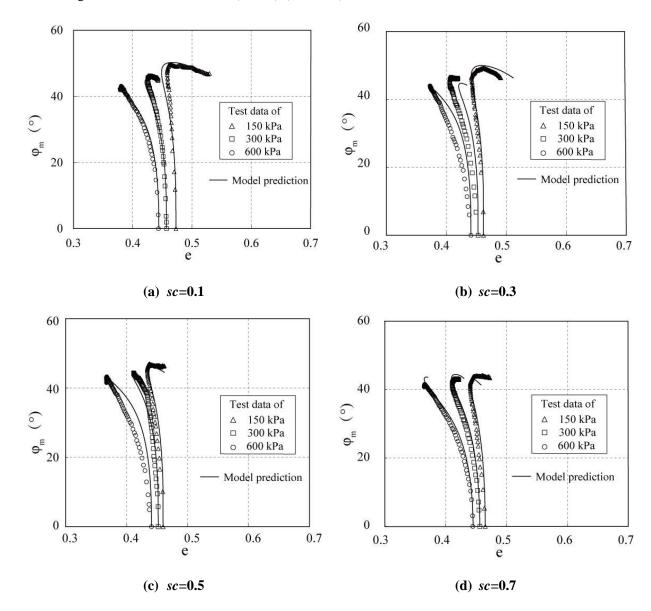
$$\sin\varphi_m = \frac{\sigma_1 - \sigma_3}{\sigma_1 + \sigma_3} \tag{21}$$

249 where  $\sigma_1$  and  $\sigma_3$  are major and minor principal stresses, respectively. In axisymmetric

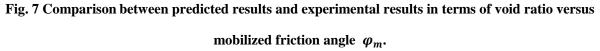
250 conditions,  $\sin\varphi_m = \frac{3\eta}{6+\eta}$ .

Figure 7 shows the comparisons between test results and model predictions in terms of the

relationship between mobilized friction angle  $\varphi_m$  and void ratio *e*. It can be seen that the model captures well the variations of  $\varphi_m$  with *e* for SRMs under different confining pressures, although the model slightly overestimates  $\varphi_m$  at high confining pressure of 600 kPa when *sc*=0.3 and 0.5. This might be due to the formation of a dual skeleton structure in the specimen with both rock and



soil grains when sc=0.3 and 0.5 (close (sc)<sub>th</sub> = 0.54).



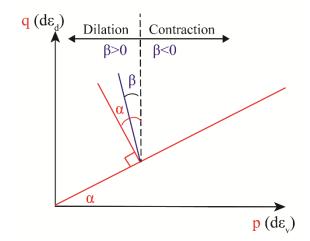
#### 257 6. Effect of soil content on essential behavior of SRMs

### 258 **6.1 Effect of** *sc* **on the non-associate behavior**

In order to investigate the effect of *sc* on the non-associated behavior of rock-dominated SRMs, the yield surface normal direction  $\alpha$  and flow rule direction  $\beta$  can be introduced as illustrated in Figure 8 in the p-q plane and dual  $d\varepsilon_{v}$ - $d\varepsilon_{d}$  plane for a non-associated flow rule with on a yield surface corresponding to Mohr-Coulomb criterion.

$$\tan \alpha = \frac{q}{p} \tag{22}$$

$$\tan\beta = -\frac{\mathrm{d}\varepsilon_{\nu}}{\mathrm{d}\varepsilon_{d}} \tag{23}$$



263

Fig. 8 Schematic diagram of normal to yield surface and flow rule direction. Both directions can be formulated with respect to angles  $\alpha$  and  $\beta$ . The sign convention of soil mechanics is

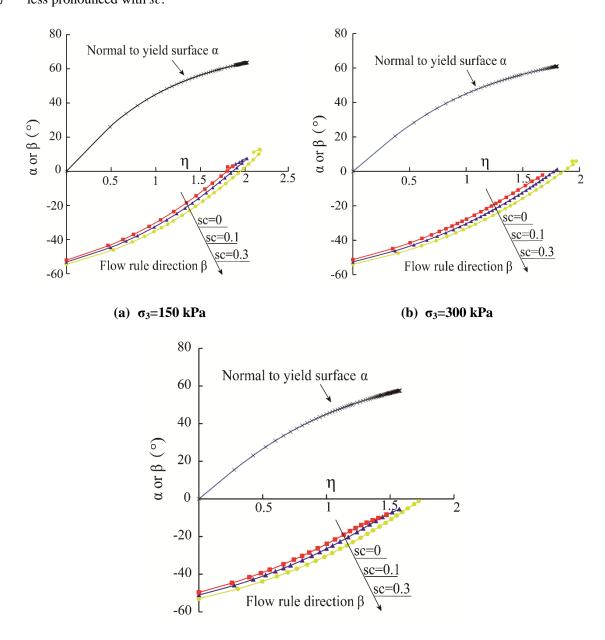
applied (positive stress in compression and positive strain in contraction).

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266

268 Knowing the model yield function from equation (12), the normal direction to the yield 269 surface  $\alpha$  and flow rule direction  $\beta$  are shown in Figure 9 for rock-dominated SRMs (*sc*=0, 0.1

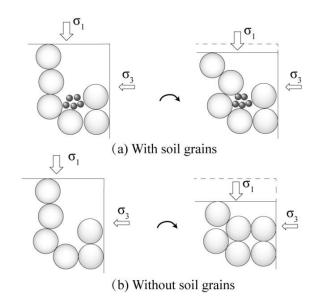
and 0.3) under confining pressures of 150 kPa, 300 kPa and 600 kPa. It can be seen that for each confining pressure, when compared at the same stress ratio  $\eta$ ,  $\alpha$  of SRMs with different *sc* are the same, indicating a non-dependence of the normal yield surface direction upon *sc*. On the other hand,  $\beta$  increases with *sc* (positive for dilation and negative for contraction), indicating that SRMs exhibit more dilantancy with the increase of *sc*. Therefore, the angle between normal direction of yield and flow rule direction decreases with *sc*, meaning the non-associate character is less pronounced with *sc*.



#### (c) σ<sub>3</sub>=600 kPa

# Fig. 9 Normal to yield surface and flow rule direction of rock-dominated SRMs with *sc*=0, 0.1 and 0.3 under confining pressure of 150 kPa, 300 kPa and 600 kPa.

277	The mesoscale origin of plastic deformation in SRMs is illustrated in Figure 10. The
278	macroscopic activation of the plastic behavior corresponds to substantial grain rearrangements
279	resulting from the collapse of preexisting force chains oriented in the principal stress direction (the
280	vertical direction in Figure 10). Once they collapse, the specimen shrinks in this direction together
281	with smaller expansion in the lateral direction. When small grains fill the pore space, the vertical
282	contraction decreases, whereas the lateral expansion is mostly unaffected. Therefore, the increase
283	in soil content is expected to enhance dilatancy in SRMs specimens.

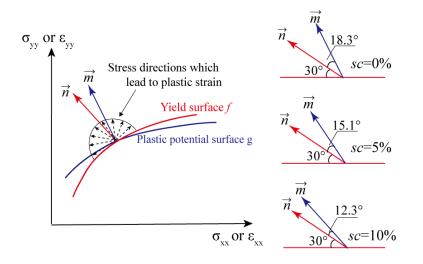


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#### Fig. 10 Illustration of the plastic deformation with and without soil grains [29]

The conceptual model of Figure 10 was proved with DEM results by Wang et al. (2021) [29] and Wautier et al. (2019) [32]. Wang et al. (2021) [29] found that for rock-dominated SRMs, *sc* does not affect the normal direction of yield surface but it changes the flow rule direction. For rock-dominated SRMs, the non-associated character is less pronounced with the increase of *sc*, as 290 shown in Figure 11. It should be noted that the normal directions of yield and plastic potential 291 surfaces shown in Fig. 11 are computed directly from DEM results, while the normal to yield 292 surface in Figure 9 depends on the expression of yield surface selected (e.g., Equation (12)). 293 Similarly, in the proposed model, soil content is found to affect the flow rule direction but the 294 mechanical state and the yield surface of rock-dominated SRMs are not affected. Indeed, this is an 295 intrinsic property of the elasto-plastic model of Li and Dafalias (2000) [9] used in this study. The 296 model assumes that the yield surface follows the stress state at any time (see Equation 12 where 297 the current stress ratio is used). Consequently, the yield surface is independent of the soil content. 298 The consistency between model predictions and DEM simulations proves the ability of this 299 constitutive model to properly reflect the non-associated properties of SRMs with sc.



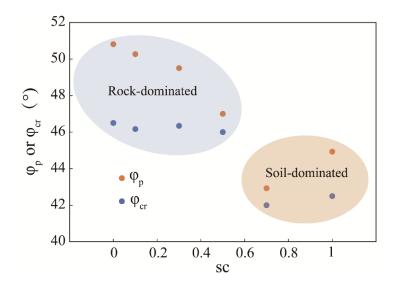
300

301 **Fig. 11** Schematic diagram of yield surface f and plastic potential surface g (on the left) and their 302 normal directions  $\vec{m}$  and  $\vec{n}$  (on the right) for DEM specimens with different *sc* at the same stress 303 ratio  $\eta$ =0.43 (Wang et al., 2021 [29]).

#### 304 **6.2 Effect of** *sc* **on peak and critical state friction angle**

305 As can be seen in Figure 7,  $\varphi_m$  increases to a peak value  $\varphi_p$ , i.e., peak friction angle, and

306 then decreases to the critical state friction angle  $\varphi_{cr}$  at relatively low confining pressures. Values 307 of  $\varphi_p$  and  $\varphi_{cr}$  of SRMs with varying sc are displayed in Figure 12. It can be seen that  $\varphi_{cr}$ 308 almost keeps constant for both rock-dominated structure (at low sc) and soil-dominated structure 309 (at high sc), while  $\varphi_p$  decreases with sc for rock-dominated structure and increases with sc for soil-dominated structure. For rock-dominated structures, as a rock skeleton exists,  $\varphi_{cr}$  is 310 primarily governed by intergranular friction between rock grains. Therefore,  $\varphi_{cr}$  is constant for 311 312 all rock-dominated SRMs. Likewise, for soil-dominated structures, a soil skeleton exists, and 313  $\varphi_{cr}$  is primarily governed by friction characteristics of soil grains. Therefore,  $\varphi_{cr}$  is also constant 314 for all soil-dominated SRMs. The reason why  $\varphi_{cr}$  in rock-dominated structures is higher than 315 that in soil-dominated structures is that the rock grains used in triaxial tests are more angular than 316 soil grains. However,  $\varphi_p$  is related to the density (in particular  $e_{sk}$ ) of the initial SRM samples. It 317 is reported in [29] that esk decreases with sc for rock-dominated SRMs and increases with sc for soil-dominated SRMs. This explains why  $\varphi_p$  decreases with sc for rock-dominated SRMs and 318 increases with sc for soil-dominated SRMs. 319



320

321

Fig. 12 Changes in peak friction angle  $\varphi_p$  and critical friction angle  $\varphi_{cr}$  with sc

322

#### 323 **7. Closure remarks**

324 The main contribution of this work is the proposal of a simple method to predict the 325 stress-strain responses of SRMs with varying soil contents. This method incorporates the concept 326 of an updated skeleton void ratio. Using this concept, SRMs with different sc should no longer be 327 regarded as different materials with their own sc-dependent model parameters. Only model 328 parameters of pure rock and of pure soil are required to describe the stress-strain response SRMs 329 with varying sc. Our proposal to adopt the skeleton void ratio is generic and can be applied to 330 other constitutive frameworks, ranging from phenomenological models to micro-mechanically based models provided that skeleton parameters (for instance,  $e_{sk}$  and  $\psi_{sk}$ ) are adopted. In this 331 332 manuscript, the constitutive framework proposed by Li and Dafalias (2000) [9] is adopted as an 333 example to demonstrate the effectiveness of this method. Extending this investigation toward 334 other constitutive frameworks will be considered in future work.

The proposed method demonstrates a satisfying ability for predicting stress-strain responses of SRMs with different soil contents and confining pressures. In addition, it successfully reflects the non-associativity behavior of rock-dominated SRMs. Some discrepancies between simulated and experimental curves are observed when *sc* is close to  $(sc)_{th}$ . One possible reason for this phenomenon is that a dual skeleton structure composed of both rock and soil grains is formed in SRM when *sc* is close to  $(sc)_{th}$ . This could be improved in future work by characterizing the complex dual skeleton structure in DEM simulations or X-ray tomography images for instance.

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#### 441 Appendix: Calibration of model parameters

- 442 (1) Elastic parameters
- 443 The initial elastic shear modulus, *G*, can be obtained from the experimental data of deviatoric
- 444 stress, q, versus the deviatoric strain,  $\varepsilon_s$  when the axial strain is lower than 0.2%. Rearrangement
- 445 of Equation (10) gives:

$$G_0 = G \frac{(1+e)}{(2.97-e)^2 \sqrt{P' P_a}}$$
(A1)

446 The values of the elastic constant,  $G_0$ , at various confining pressures, can be determined from

Equation (A1). The average value of  $G_0$  under different confining pressures is adopted.

448 Based on Equations (7)-(11), the Poisson's ratio,  $\nu$  can be obtained by:

$$v = \frac{9d\varepsilon_s^e - 2d\varepsilon_v^e}{18d\varepsilon_s^e + 2d\varepsilon_v^e} \approx \frac{9\varepsilon_s - 2\varepsilon_v}{18\varepsilon_s + 2\varepsilon_s}$$
(A2)

#### 449 (2) Critical state parameters

450  $e_{\Gamma}$ ,  $\lambda$  and  $\xi$  can be determined by directly fitting the experimental data for the critical state 451 line. The critical state stress ratio *M* can be obtained by fitting critical state test data in *p'-q* plane 452 with a function of q=Mp'.

#### 453 (3) Dilatancy parameters

454 The parameter *m* is determined from Equation (17) at a phase transformation state, at which 455 D=0, and thus,

$$m = \frac{1}{\psi^d} ln \frac{M^d}{M} \tag{A3}$$

456 where  $\psi^d$  and  $M^d$  are the values of  $\psi$  and  $\eta$  at the phase transformation state.

457 Ignoring the small elastic strain, we have:

$$\frac{\mathrm{d}\varepsilon_{\nu}}{\mathrm{d}\varepsilon_{q}} \approx \frac{\mathrm{d}\varepsilon_{\nu}^{p}}{\mathrm{d}\varepsilon_{q}^{p}} = D = d_{0} \left( \exp(m\psi) - \frac{\eta}{M} \right) \tag{A4}$$

458 The parameter  $d_0$  is determined based on the  $d\varepsilon_v - d\varepsilon_q$  curve.

#### 459 (4) Hardening parameters

460 The parameter *n* is determined by Equation (14) at a peak stress state, at which  $K_p=0$ :

$$n = \frac{1}{\psi^b} ln \frac{M}{M^b} \tag{A5}$$

461 where  $\psi^b$  and  $M^b$  are the values of  $\psi$  and  $\eta$  at the peak stress state.

462 Combining Equations (10) (13) and (14) for conventional drained tests (dp'=dq/3) yields:

$$\frac{\mathrm{d}q}{\mathrm{d}\varepsilon_q} \approx \frac{\mathrm{d}q}{\mathrm{d}\varepsilon_q^p} = \frac{K_p}{1 - \eta/3} = h \left\{ \frac{G_0 (2.97 - e)^2 \sqrt{p' p_a} \left(\frac{M}{\eta} - \exp(n\psi)\right)}{(1 + e)(1 - \eta/3)} \right\}$$
(A6)

463 As all the model parameters in the brackets are known, *h* is determined based on  $dq - d\varepsilon_q$ 464 curves along drained triaxial loading paths. Then parameters  $h_1$  and  $h_2$  can be obtained by 465 equation  $h = h_1 - h_2 e_0$ .