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To cite this version:

Thomas Guilmeau, Alain Rapaport. Multiplicity of periodic orbits with coexistence in the chemostat subject to periodic removal rate. 2023. hal-03982233v1

HAL Id: hal-03982233
https://hal.inrae.fr/hal-03982233v1

Preprint submitted on 10 Feb 2023 (v1), last revised 20 Oct 2023 (v2)

Multiplicity of periodic orbits with coexistence in the chemostat subject to periodic removal rate

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February 9, 2023

Abstract

We address the currently open problem of existence of multiple periodic orbits in the chemostat model with periodic removal rate. We give conditions on a subset of growth functions that ensure the coexistence of an arbitrary number of species within this subset. We show that proportions of some powers of the species densities are periodic functions, leading to an infinity of distinct periodic orbits depending on the initial condition. We give also conditions allowing the coexistence of two distinct subsets of species. Although these conditions are non-generic, we show in simulations that when these conditions are only approximately satisfied, then the behavior of the solutions are close from the non-generic case over a long time interval, justifying the interest of our study.

Key-words. Chemostat model, Periodic removal rate, Coexistence, Poincaré map, Multiplicity of periodic orbits.

AMS Subject Classification. 34C25, 37C25, 92-10, 92D25, 92D40.

1 Introduction

The mathematical model of the chemostat can represent a vast array of natural phenomena where different living species compete for a common limiting resource. In particular, it is widely used to model waste water bioprocesses or to capture competition for a single resource in ecological modeling [15, 5]. Recall that the chemostat is originally an experimental device in which bacterial species grow in a perfectly stirred vessel of constant volume, continuously removed and fed with fresh substrate. Of course, the mathematical model of the chemostat reflects also the behavior of this device.

When several species (of densities x_i, $i = 1, \cdots$) all consume a single substrate that is fed in the ecosystem at a constant rate, then the mathematical theory of the chemostat claims that the *Competitive Exclusion Principle* holds, meaning that all the species will asymptotically disappear from the system, except the fittest one [15, 5]. However, in more complex environments, this property does not always hold. This is for instance the case of periodic operation of bioreactors [13] or temporal fluctuations in ecology [6, 4, 10]. If the environment favors for some time one species and then another one in a balanced way, then one may expect the two species to coexist. A large part of the literature is dedicated to the study of the asymptotic behavior of theses dynamics, with two or more species, under periodic removal rate [16, 1, 7, 9] or periodic nutrient input [6, 14, 4, 17] or both [18, 11].

Indeed, when a vessel with two species of concentrations $x_i(t)$ at time t ($i = 1, 2$) is fed with a periodic removal rate, the system can exhibit the following different behaviors :

- The two species are washed out, that is $\lim_{t\to+\infty} x_i(t) = 0$, for $i = 1, 2$.
- One species only is washed out while the other survives.
- The two species coexist, that is $x_i(t) > \epsilon > 0$, for $t \geq 0$ and $i = 1, 2$, for some $\epsilon > 0$.

Those behaviors can be predicted using the Floquet theory. Namely, the Floquet exponents allow to determine if the periodic solutions with only one species are unstable. If both of these single-species periodic solutions are unstable, then the theory of competitive planar systems [2] applied to the chemostat model shows that any positive solution converges asymptotically to a periodic solution with species coexistence (a complete description of this theory is given for instance in [15]). However, the possibility of having several attracting periodic orbits remains an open problem. Although there is no theoretical obstruction for this, no such example has been yet exhibited in the literature. The purpose of the present work is to show that having multiplicity of periodic orbits with coexistence is indeed possible, but for a particular class of growth functions, providing moreover an infinite number of periodic orbits, with two or more species.

The construction that we propose here is non-generic, relying on a relatively strong condition on the growth functions that is deemed as unfeasible in real life. However, one may face practical situations close from this non-generic case, which may provide a better understanding of the periodic chemostat over a long duration. Moreover, our construction may be leveraged to design sufficient conditions for uniqueness of periodic orbits with coexistence.

Our construction lies on the concept of "taxon", which in this paper, denotes a group of species whose growth functions share the same shape. We show that the fittest species from each taxon behave like one species in some sense, and when they can coexist, then there exists an infinite number of periodic orbits. Moreover, when two taxa are present in the chemostat, we show that the fittest species from each taxon can coexist all together under conditions that generalize the coexistence conditions of two species in the chemostat with periodic removal rate.

The paper is organized as follows. In the next Section 2, we recall the equations of the model with the usual assumptions and give some useful definitions and preliminary results. In Section 3, we present our new conditions with a concept of "taxon" and show some relevant properties. Sections 4 and 5 give our main results, first for a single taxon and then for competition between taxa. Finally, we presents and discusses several numerical simulations to illustrate our results Section 6, before drawing conclusions in Section 7.

2 The setting

We consider the multi-species chemostat model written as follows

$$\begin{aligned} \dot{s} &= u(t)(s_{in}-s) - \sum_{i=1}^{n} \frac{\mu_i(s)}{Y_i} x_i, \\ \dot{x}_i &= (\mu_i(s) - u(t))x_i, \quad 1 \leq i \leq n, \end{aligned} \tag{1}$$

(with $n \geq 2$) where s is the substrate concentration, s_{in} the input substrate concentration and x_i ($1 \leq i \leq n$) are the respective concentrations of the n populations. We recall the well-known fact that the yield coefficients Y_i, $1 \leq i \leq n$ can be taken equal to one without loss of generality, by a change of variables (x_i/Y_i replaced by x_i). As usual, the growth functions $\mu_i(.)$ satisfy the following properties.

Assumption 1. *For any $i = 1, \cdots, n$, μ_i is a C^1 increasing function from $\mathbb{R}_+$ to $\mathbb{R}_+$ with $\mu_i(0) = 0$.*

Under this assumption, we define classically the *break-even concentration* for each species $i = 1, \cdots, n$ as the function

$$\lambda_i(v) := \sup \{ s \in \mathbb{R}_+ ; \ \mu(s) < v \}, \quad v > 0.$$

Note that when $\lambda_i(v)$ is finite, one has necessarily $\mu_i(\lambda_i(v)) = v$.

Assumption 2. *The removal rate $u(\cdot)$ is a time measurable function from $[0, +\infty)$ to $[u_-, u_+]$ with $0 < u_- \leq u_+$, which is T-periodic (with $0 < T < +\infty$). We posit*

$$\bar{u} = \frac{1}{T} \int_0^T u(t)\, dt.$$

Consider now the variable $b := s + \sum_{i=1}^{n} x_i$, whose dynamics is given by

$$\dot{b} = u(t)(s_{in} - b). \tag{2}$$

Under Assumption 2, one has $u(t) \geq u_- > 0$ at any $t \geq 0$ and one can deduce that the solution of (2) converges exponentially to s_{in}, independently of the initial condition of (1). With Assumption 1, the solutions of (1) are uniquely defined and bounded for any non-negative initial condition. Therefore the asymptotic behavior of the $n+1$ dimensional system (1) is determined by the n dimensional dynamics

$$\dot{x}_i = \left[\mu_i\left(s_{in} - \sum_{j=1}^{n} x_j\right) - u(t)\right] x_i, \quad 1 \leq i \leq n, \tag{3}$$

which leaves the set

$$\Delta_0 := \left\{x \in \mathbb{R}_+^n, \text{ such that } \sum_{i=1}^{n} x_i \leq s_{in}\right\}$$

forwardly invariant. In the remaining, we shall consider the asymptotic dynamic (3) on the set Δ_0 only (which is biology relevant in the chemostat framework).

Note that solutions of (3) with $x_i(0) = 0$ for some i are such that $x_i(t) = 0$ for any t. Therefore, by uniqueness of solutions of (3), we deduce that a solution $x(\cdot)$ of (3) with a positive initial condition has to stay positive for any time.

We give below a Lemma that will be useful in the rest of the paper.

Lemma 1. *There exists a number* $\underline{s} \in (0, s_{in})$ *such that the subset*

$$\Delta := \left\{x \in \Delta_0\,;\, s_{in} - \sum_{i=1}^{n} x_i > \underline{s}\right\}$$

is forward invariant and attractive by the dynamics (3)*.*

Proof. Posit $s = s_{in} - \sum_{i=1}^{n} x_i$. Since x is a solution of (3) and the set Δ_0 is forward invariant, it comes

$$\begin{aligned}
\dot{s} &= -\sum_{i=1}^{n} \mu_i(s) x_i(t) + u(t)(s_{in} - s) \\
&\geq -\left(\max_{1 \leq i \leq n} \mu_i(s)\right) \sum_{i=1}^{n} x_i(t) + u(t)(s_{in} - s) \\
&\geq -\left(\max_{1 \leq i \leq n} \mu_i(s)\right) s_{in} + u(t)(s_{in} - s) \\
&\geq g(s) := -\left(\max_{1 \leq i \leq n} \mu_i(s)\right) s_{in} + u_-(s_{in} - s).
\end{aligned}$$

The function g is continuous and decreasing with $g(0) = u_- s_{in} > 0$ and $g(s_{in}) < 0$. By the intermediate value Theorem, there exists a number $s_m \in (0, s_{in})$ such that $g(s_m) = 0$ with $g(s) > 0$ for $s < s_m$. Therefore, for any $\underline{s} \in (0, s_m)$, the domain $\{s > \underline{s}\}$ is forwardly invariant and attractive by the dynamics of s, which amounts to claim that the subset Δ is forward invariant and attractive by the dynamics (3). ☐

3 A taxonomic assumption and its consequences

We shall consider subsets $I \subset \{1, \cdots, n\}$ of at least two species, whose growth functions μ_i ($i \in I$) share a common property, defining what we propose to call a *taxon* in the present context.

Definition 1. *A subset of populations* $I \subset \{1, \cdots, n\}$ *belong to a same* taxon *if there exists a* C^1 *increasing function* φ *with* $\varphi(\underline{s}) < u_-$ *and* $\varphi(s_{in}) > u_+$*, and numbers* $\alpha_i > 0$*,* $\beta_i \in \mathbb{R}$ *for* $i \in I$ *such that*

$$\mu_i(s) = \alpha_i \varphi(s) + \beta_i, \quad s \in [\underline{s}, s_{in}], \quad i \in I. \tag{4}$$

We shall say that such a function φ *is a generating growth function of the taxon. We also define the subset* $I^\star \subset I$*, which is such that*

$$I^\star := \{i \in I;\ \lambda_i(\bar{u}) = \min_{j \in I} \lambda_j(\bar{u})\}.$$

In this definition, the choice of the generating function φ is not unique but it can typically represent a *canonical growth function* that verifies Assumption 1 so that the growth functions μ_i among a taxon differ only by an affine transformation away for 0 i.e. on the interval $[\underline{s}, s_{in}]$. Note that condition (4) cannot be imposed for any $s > 0$ if $\beta_i \neq 0$, because growth functions have to be equal to 0 at the origin (Assumption 1). Several examples of growth functions that satisfy Assumption 1 and condition (4) will be given in Section 6. Let us point out that this condition does not prevent the graphs of the functions μ_i to cross on the domain $[\underline{s}, s_{in}]$.

The second part of Definition 1 concerns the subset $I^\star \subset I$, which is made of all the species that share the same minimal break-even concentration. In some sense, the set $I^\star$ contains the fittest species, as these are the surviving species in the autonomous chemostat model [5]. Generically, the set $I^\star$ is reduced to a singleton, but we shall study in this work the non-generic situation where more than one species belong to $I^\star$.

For convenience, we shall define some auxiliary numbers. The function φ being increasing with $\varphi(\underline{s}) < u_-$ and $\varphi(s_{in}) > u_+$, and as $\bar{u} \in [u_-, u_+]$, there exists an unique $\bar{s} \in (\underline{s}, s_{in})$ such that

$$\varphi(\bar{s}) = \bar{u}, \tag{5}$$

and for any $i \in I$ we define the numbers

$$\gamma_i := \alpha_i^{-1}(\mu_i(\bar{s}) - \bar{u}) = \frac{\beta_i}{\alpha_i} + \bar{u}\left(1 - \alpha_i^{-1}\right). \tag{6}$$

Remark 1. *We have for any $i \in I$ that $\mu_i(\lambda_i(\bar{u})) = \bar{u}$, which is equivalent to write*

$$\varphi(\lambda_i(\bar{u})) = \alpha_i^{-1}(\bar{u} - \beta_i) = \bar{u} - \gamma_i, = \varphi(\bar{s}) - \gamma_i.$$

Since the generating growth function φ is increasing, a higher value γ_i corresponds to a lower break-even concentration λ_i. In particular, we have that $i \in I^\star$ exactly when $\gamma_i = \max_{j \in I} \gamma_j$.

The key point in our study will be to consider particular ratios of powers of species concentrations, for species that belong to I, defined as follows

$$\rho_{ij} := \frac{\sqrt[\alpha_i]{x_i}}{\sqrt[\alpha_j]{x_j}}, \quad i, j \in I.$$

The dynamics of these variables present some particular properties, as shown in the next Proposition.

Proposition 1. *Let I be a subset of populations that belong to a same taxon. For any solution positive solution in Δ, the dynamics of the ratios ρ_{ij} are as follows*

$$\dot{\rho}_{ij} = \Big(\gamma_i - \gamma_j + (\bar{u} - u(t))(\alpha_i^{-1} - \alpha_j^{-1})\Big)\rho_{ij}, \quad i, j \in I. \tag{7}$$

Proof. Let us first differentiate the equality $\sqrt[\alpha_j]{x_j}\rho_{ij} = \sqrt[\alpha_i]{x_i}$ with respect to t:

$$\frac{1}{\alpha_j} x_j^{\frac{1}{\alpha_j} - 1} \dot{x}_j \rho_{ij} + x_j^{\frac{1}{\alpha_j}} \dot{\rho}_{ij} = \frac{1}{\alpha_i} x_i^{\frac{1}{\alpha_i} - 1} \dot{x}_i, \quad i, j \in I$$

and replace $\dot{x}_k$ for $k = i, j$ by the expression $(\alpha_k \varphi(s(t)) + \beta_k - u(t))x_k$. One obtains

$$x_j^{\frac{1}{\alpha_j}} \left(\varphi(s(t)) + \frac{\beta_j}{\alpha_j} - \alpha_j^{-1} u(t)\right) \rho_{ij} + x_j^{\frac{1}{\alpha_j}} \dot{\rho}_{ij} = x_i^{\frac{1}{\alpha_i}} \left(\varphi(s(t)) + \frac{\beta_i}{\alpha_i} - \alpha_i^{-1} u(t)\right), \quad i, j \in I.$$

Multiplying by $x_j^{\alpha_j}$, one can write

$$\dot{\rho}_{ij} = \left(\frac{\beta_i}{\alpha_i} - \frac{\beta_j}{\alpha_i} - u(t)(\alpha_i^{-1} - \alpha_j^{1})\right) \rho_{ij}, \quad i, j \in I.$$

Finally, from the definition (6) of numbers γ_k, $k \in i, j$, one can also write

$$\dot{\rho}_{ij} = \Big(\gamma_i - \gamma_j + (\bar{u} - u(t))(\alpha_i^{-1} - \alpha_j^{-1})\Big)\rho_{ij}, \quad i, j \in I.$$

□

The dynamics (7) of the ratios ρ_{ij} within a taxon presents thus the remarkable feature that the time evolution of each ratio depends only on its initial value and the function u, i.e. their dynamics are decoupled. As a consequence, one obtains the following properties of the solutions of (3).

Proposition 2. *Let I be a subset of populations that belongs to a same taxon. For any positive solution in Δ, one has*

1. *For i, j in $I^\star$, ρ_{ij} are periodic functions.*
2. *For $i \in I \setminus I^\star$, x_i converges asymptotically to 0.*
3. *If $\liminf_{t>0} x_i(t) > 0$ for some $i \in I^\star$, then $\liminf_{t>0} x_j(t) > 0$ for any other $j \in I^\star$.*

Proof. From (7), one obtains the expression

$$\frac{d}{dt}\log \rho_{ij} = (\gamma_i - \gamma_j) + (\bar{u} - u(t))(\alpha_i^{-1} - \alpha_j^{-1}), \quad t \geq 0$$

that we integrate between t and $t+T$:

$$\log \rho_{ij}(t+T) = \log \rho_{ij}(t+T) + (\gamma_i - \gamma_j)T, \quad t \geq 0$$

which gives equivalently

$$\rho_{ij}(t+T) = \rho_{ij}(t)e^{(\gamma_i - \gamma_j)T}, \quad t \geq 0.$$

Then, for $i, j \in I^\star$, $\gamma_i = \gamma_j$ and the function ρ_{ij} is thus periodic. If $i \notin I^\star$, for any $j \in I^\star$, one has $\gamma_j > \gamma_i$ (see Remark 1) and thus $x_i(t) \to 0$ for $t \to +\infty$. The last point of the Proposition is a straightforward consequence of point 1. ☐

This result states that the Competitive Exclusion Principle occurs within a taxon in the periodic chemostat. It also means that when one or several species of a same taxon persist in a periodic chemostat, it can be invaded by a new one belonging to the same taxon preserving the coexistence of all resident populations, under the condition that all species have the same minimal break even concentration (for the average removal rate). Diversity can be then (theoretically) augmented within a same taxon.

The (non-generic) property of having identical break-even concentrations that implies coexistence of species is already known in the classical chemostat model with constant removal rate (see for instance [5]). However, it is also known that this condition does not guarantee the coexistence under periodic removal rate. Instead, integral conditions which depends on the periodic function u have to be fulfilled [15]. Here, the remarkable feature within a taxon is that under the simple condition of equal break-even concentrations, coexistence can be guaranteed whatever is the periodic function u (provided that its average value is equal to the fixed value $\bar{u}$). Moreover, we show below that a multiplicity of periodic solutions can be obtained within a taxon, which is a new result in the theory of the periodic chemostat model, up to our knowledge.

Now and for the rest of the paper, we shall assume that each species population can persist alone, which is ensured by the following hypothesis.

Assumption 3. *One has $\mu_i(s_{in}) > \bar{u}$ for any $i = 1, \cdots, n$.*

4 Multiplicity of periodic solutions within a single taxon

We consider here that all the species belong to a same taxon.

Proposition 3. *Assume that the whole set $I = \{1, \cdots, n\}$ of populations belong to a same taxon. Then for any positive initial condition in Δ, the solution of* (3) *converges asymptotically to a periodic solution composed of all species in $I^\star$, the other species being washout. Moreover, if $I^\star$ is not reduced to a singleton, the system* (3) *admits an infinite number of periodic solutions in Δ with distinct orbits, which depend on the initial proportions of species in $I^\star$ only.*

Proof. Let $x(\cdot)$ be a positive solution of (3) in Δ and $\rho_{ij}(\cdot)$, $i,j \in I$, be the corresponding solutions of (7). Take $i^\star$ in $I^\star$ such that

$$\alpha_{i^\star} = \min_{i \in I^\star} \alpha_i.$$

We show that $x_{i^\star}$ converges asymptotically to a positive periodic solution.

The variable $x_{i^\star}$ can be written as the solution of the non-autonomous scalar dynamics

$$\dot{x}_{i^\star} = (f(t, x_{i^\star}) - u(t))x_{i^\star}$$

where

$$f(t,y) := \mu_{i^\star}\left(s_{in} - \sum_{i \in I^\star} \rho_{ii^\star}(t)^{\alpha_i}\, y^{\frac{\alpha_i}{\alpha_{i^\star}}} - \sum_{i \notin I^\star} x_i(t)\right).$$

Let $\tilde{\mu}_{i^\star}$ be a C^1 increasing extension of the function $\mu_{i^\star}$ for negative arguments, and consider the dynamics

$$\dot{y} = (\tilde{f}(t,y) - u(t))y \tag{8}$$

on $\mathbb{R}$, where $\tilde{f}$ consists in replacing $\mu_{i^\star}$ by $\tilde{\mu}_{i^\star}$ in the expression of the function f. Note that the choice of $x_{i^\star}$ implies that one has $\alpha_i/\alpha_{i^\star} \geq 1$ for any $i \in I^\star$. Therefore, $\tilde{f}$ is Lipschitz with respect to y and the solutions of (8) are well defined. At $y = s_{in}$, the argument of $\tilde{\mu}_{i^\star}$ is negative, but as $\mu_i(0) = 0$ and $\tilde{\mu}_{i^\star}$ is increasing, one has necessarily $\tilde{\mu}_{i^\star} < 0$ for negative arguments, and then $\tilde{f}(t, s_{in}) \leq 0$ for any $t \geq 0$. The set $[0, s_{in}]$ is thus forwardly invariant. Clearly, $x_{i^\star}$ is the solution of (8) for the initial value $y(0) = x_{i^\star}(0)$ (which belongs to $[0, s_{in}]$).

According to Proposition 2, the functions $\rho_{ii^\star}$ for $i \in I^\star$ are T-periodic and variables x_i for $i \notin I^\star$ converge asymptotically to 0. We then consider the limiting function

$$\tilde{f}_l(t,y) := \tilde{\mu}_{i^\star}\left(s_{in} - \sum_{i \in I^\star} \rho_{ii^\star}(t)^{\alpha_i}\, y^{\frac{\alpha_i}{\alpha_{i^\star}}}\right)$$

which is T-periodic and verifies

$$\lim_{t \to +\infty} |\tilde{f}(t,y)y - \tilde{f}_l(t,y)y| = 0 \text{ uniformly for } y \in [0, s_{in}].$$

Consequently, by Proposition 3.2 in [19], the non-autonomous semi-flow of (8) in $[0, s_{in}]$ is asymptotically periodic with limit periodic semi-flow of

$$\dot{y} = (\tilde{f}_l(t,y) - u(t))y \tag{9}$$

(for which $[0, s_{in}]$ is also forwardly invariant).

We follow now the approach exposed in [15] for one dimensional periodic dynamics, but adapted here to our context. Let us consider the Poincaré map P associated to the periodic dynamics (9)

$$P : y_0 \in [0, s_{in}] \mapsto y(T, y_0) \in [0, s_{in}]$$

where $y(\cdot, y_0)$ denotes the solution of (9) with $y(0) = y_0$. One has clearly $P(0) = 0$, and from the Theorem of continuous dependency of the solution of ordinary differential equation with respect to the initial condition, P is continuously differentiable with $P'(y_0) = z(T)$, where $z(\cdot)$ is solution of

$$\dot{z} = \Big(\partial_y \tilde{f}_l(t, y(t,y_0))y(t,y_0) + \tilde{f}_l(t, y(t,y_0)) - u(t)\Big)z, \quad z(0) = 1$$

that is

$$z(T) = \exp\left(\int_0^T \partial_y \tilde{f}_l(t, y(t,y_0))y(t,y_0) + \tilde{f}_l(t, y(t,y_0)) - u(t)\, dt\right) > 0.$$

The map P is thus increasing and one has $P'(0) = \exp\big(T\left(\mu_{i^\star}(s_{in}) - \bar{u}\right)\big) > 1$ (by Assumption 3). So 0 is a repulsive fixed point of the map P, and for any $y_0 > 0$, the sequence $\left\{P^k y_0\right\}_{k \in \mathbb{N}}$ is strictly monotonic and bounded, thus converging to a positive fixed point $y^\star$ of P. Moreover, as $y(\cdot, y^\star)$ is periodic, one has

$$\int_0^T \tilde{f}_l(t, y(t, y^\star)) - u(t)\, dt = 0$$

and thus

$$P'(\bar{y}) = \exp\left(\int_0^T \partial_y \tilde{f}_l(t, y^\star) y(t, y^\star) dt\right).$$

As the functions $\tilde{\mu}_i$ are assumed to be increasing, one has $\tilde{\mu}'_{i^\star} > 0$ which implies $\partial_y \tilde{f}_l < 0$, and thus $P'(\bar{y}) < 1$. Therefore the map $y \mapsto P(y) - y$ is decreasing at each root, which implies that it cannot have more than one root. We conclude that the (positive) fixed point $y^\star$ is unique.

As P admits a finite number of fixed points (indeed only one), one can apply the results about asymptotically autonomous discrete dynamical systems (Theorem 2.4 in [19]), from which one gets

$$\lim_{k \to +\infty} x_{i^\star}(t + kT) = y(t, y^\star), \quad t \in [0, T].$$

and we conclude that the solution $x(\cdot)$ converges asymptotically to a periodic solution of (3) in Δ, for which all species in $I^\star$ are present, the other species being excluded:

$$\lim_{k \to +\infty} |x(t + kT) - x^p(t)| = 0, \quad t \in [0, T]$$

where $x^p(\cdot)$ is the periodic solution given by

$$x_i^p(t) = \left\{ \begin{array}{ll} \rho_{ii^\star}(t)^{\alpha_i}\, y(t, y^\star)^{\frac{\alpha_i}{\alpha_{i^\star}}} > 0, & i \in I^\star, \\ 0, & i \notin I^\star, \end{array} \right.$$

for $t \in [0, T]$.

Consider now another positive initial condition in Δ but with the same initial ratios $\rho_{ii^\star}(0)$ for $i \in I^\star$. According to (7), the functions $\rho_{ii^\star}$ ($i \in I^\star$) are identical and consequently the limiting periodic dynamics (9) is also identical. As this later one admits an unique periodic solution, we conclude that the solution of (3) converges asymptotically to the same periodic solution $x^p(\cdot)$.

We now show how to construct an infinity of distinct periodic solutions, when $I^\star$ is not reduced to a singleton. Consider a sequence $\{\rho_0^k\}_{k \in \mathbb{N}}$ of positive vectors in $\mathbb{R}^n$ such that

$$\max_{i \in I^\star \setminus \{i^\star\}} (\rho_0^{k+1})_i > \max_{t \in [0,T]} \max_{i \in I^\star \setminus \{i^\star\}} \rho_{ii^\star}^k(t), \quad k \in \mathbb{N} \tag{10}$$

where $\rho_{ii^\star}^k(\cdot)$ are the periodic solutions of (7) with $\rho_{ii^\star}^k(0) = (\rho_0^k)_i$ for $i \in I^\star \setminus \{i^\star\}$. Condition (10) imposes that for each $i \in I^\star \setminus \{i^\star\}$, the orbits $\gamma^+(\rho_{ii^\star}^k) = \{\rho_{ii^\star}^k(t),\ t \in [0, T]\}$ of (7) are all disjoint for $k \in \mathbb{N}$. Moreover, for each k, there exists an unique periodic solution $x^{pk}(\cdot)$ of (3) in Δ for which all species in $I^\star$ are present with ratios given by the functions $\rho_{ii^\star}^k$. This implies that the periodic orbits $\gamma^+(x^{pk}) = \{x^{pk}(t),\ t \in [0, T]\}$, $k \in \mathbb{N}$, of (3) are all disjoint. Indeed, if $\gamma^+(x^{pk}) = \gamma^+(x^{pl})$ for some $k \neq l$, there should exist $\tau \geq 0$ such that $x^{pk}(t) = x^{pl}(t + \tau)$ for any $t \geq 0$. In particular, one should have $x_i^{pk}(t) = x_i^{pl}(t + \tau)$ for $i \neq i^\star$ in $I^\star$ and any $t \geq 0$, but as the orbits $\gamma^+(\rho_{ii^\star}^k)$, $\gamma^+(\rho_{ii^\star}^l)$ are disjoint, one should have $\rho_{ii^\star}^k(t+\tau) \neq \rho_{ii^\star}^l(t)$ for some t, that is $x_{i^\star}^{pk}(t) \neq x_{i^\star}^{pl}(t+\tau)$ and thus a contradiction with with the fact that the orbits are non distinct. ☐

The results of Proposition 3 are twofold. First, they complete those of Proposition 2, since we now have the persistence of the species of $I^\star$, while the less fit species of $I \setminus I^\star$ are washed-out. This stronger form of the competitive exclusion principle comes from Assumption 3 and the fact that there is no species outside the taxon I. This result allows coexistence in the periodic setting under the non-generic assumption that the surviving species belong to $I^\star$. It generalizes known similar results in the non-autonomous setting [5].

Then, the second aspect of this result concerns the number of distinct periodic coexistence solutions of the system (3), which is shown to be infinite under our assumptions. While the standard theory of the periodic chemostat gives sufficient conditions for the existence of such solutions [15], it says nothing about the uniqueness of such solutions. Our construction sheds new lights on this issue, since Proposition 3 shows that if the growth functions are *close enough* to each other, in the sense that the corresponding species belong to $I^\star$, then there are infinitely many periodic coexistence solutions. We may expect that eventual sufficient conditions for uniqueness forbid growth functions to be too close in a sense close to ours.

5 Multiplicity of periodic solutions with more than one taxon

We show here that it is possible to have coexistence of two taxa in competition, each of them preserving the proportions of species having the same break-even concentrations, leading to an infinite number of periodic orbits. For technicalities, we need in this section the following additional hypothesis.

Assumption 4. *The functions μ_i, $i \in \{1, \cdots, n\}$, are analytic at any $s > \underline{s}$, and u is an analytic function of t.*

For convenience, we shall denote for any integrable scalar function $\zeta(\cdot)$ the average quantity by

$$\langle \zeta \rangle_T := \frac{1}{T} \int_0^T \zeta(t) dt.$$

Proposition 4. *Assume that one has $\{1, \cdots, n\} = I_a \sqcup I_b$, where species in I_a, resp. I_b, belong to a same taxon. For any fixed positive initial condition in Δ, let $(s_a^p(\cdot), x_a^p(\cdot))$, resp. $(s_b^p(\cdot), x_b^p(\cdot))$ be the asymptotic periodic solution of* (1) *when only species in I_a, resp. I_b, are initially present. If the conditions*

$$\lambda_{ba} := \langle \mu_i(s_b^p) \rangle_T - \bar{u} > 0, \quad i \in I_a^\star, \tag{11}$$

$$\lambda_{ab} := \langle \mu_i(s_a^p) \rangle_T - \bar{u} > 0, \quad i \in I_b^\star, \tag{12}$$

are fulfilled, then the solution of (3) *converges asymptotically to a periodic solution for which all species in $I_a^\star$ and $I_b^\star$ are present, the other species being washed out. Moreover, when there exists at least one positive initial condition satisfying the above conditions, and at least one of the subsets $I_a^\star$, $I_b^\star$ is not reduced to a singleton, there exists an infinity of distinct periodic orbits of* (3) *for which all the species in $I_a^\star \sqcup I_b^\star$ are present.*

Proof. The taxa are characterized by generating functions φ_a, φ_b and numbers $\alpha_i > 0$, $\beta_i \geq 0$ such that

$$\mu_i(s) = \alpha_i \varphi_a(s) + \beta_i, \ i \in I_a, \quad \mu_i(s) = \alpha_i \varphi_b(s) + \beta_i, \ i \in I_b.$$

Take $i_a^\star \in I_a^\star$ and $i_b^\star \in I_b^\star$. Let $x(\cdot)$ be a positive solution of (3) in Δ and define the functions, for $i \in I_a^\star \sqcup I_b^\star$

$$r_i(t) = \begin{cases} \rho_{ii_a^\star}(t), & i \in I_a^\star, \\ \rho_{ii_b^\star}(t), & i \in I_b^\star, \end{cases} \qquad t \geq 0. \tag{13}$$

Then, variables $x_{i_a^\star}$, $x_{i_b^\star}$ are solutions of the non-autonomous planar dynamics

$$\begin{aligned} \dot{y}_a &= (f_a(t, y_a, y_b) - u(t)) y_a, \\ \dot{y}_b &= (f_b(t, y_a, y_b) - u(t)) y_b, \end{aligned}$$

with

$$f_a(t, y_a, y_b) := \mu_{i_a^\star}(s(t, y_a, y_b)), \quad f_b(t, y_a, y_b) := \mu_{i_b^\star}(s(t, y_a, y_b)),$$

where

$$s(t, y_a, y_b) = s_{in} - \sum_{i \in I_a^\star} r_i(t)^{\alpha_i} y_a^{\frac{\alpha_i}{\alpha_{i_a^\star}}} - \sum_{i \in I_b^\star} r_i(t)^{\alpha_i} y_b^{\frac{\alpha_i}{\alpha_{i_b^\star}}} - \sum_{i \notin I_a^\star \sqcup I_b^\star} x_i(t).$$

Let $\tilde{\mu}_{i_a^\star}$, $\tilde{\mu}_{i_b^\star}$ be C^1 increasing extensions of the functions $\mu_{i_a^\star}$, $\mu_{i_b^\star}$ for negative arguments, and consider the dynamics

$$\begin{aligned} \dot{y}_a &= (\tilde{f}_a(t, y_a, y_b) - u(t)) y_a, \\ \dot{y}_b &= (\tilde{f}_b(t, y_a, y_b) - u(t)) y_b, \end{aligned} \tag{14}$$

in the plane, where the functions $\tilde{f}_a$, $\tilde{f}_b$ are defined with the expressions of f_a, f_b replacing the functions $\mu_{i_a^\star}$, $\mu_{i_b^\star}$ by their extensions $\tilde{\mu}_{i_a^\star}$, $\tilde{\mu}_{i_b^\star}$. As one has $\alpha_i / \alpha_{i_a^\star} \geq 1$ for $i \in I_a^\star$, and $\alpha_i / \alpha_{i_b^\star} \geq 1$ for $i \in I_b^\star$, this dynamics is Lipschitz in (y_a, y_b). Moreover the set

$$\mathcal{S} := [0, s_{in}] \times [0, s_{in}]$$

is forwardly invariant as one has $f_a(t, s_{in}, y_b) < 0$, $f_b(t, y_a, s_{in}) < 0$ for any $(y_a, y_b) \in \mathcal{S}$ and $t \geq 0$. Solutions of (14) are thus well defined in $\mathcal{S}$ and unique. The pair $(x_{i_a^\star}(\cdot), x_{i_b^\star}(\cdot))$ is such a solution for $(y_a(0), y_b(0)) = (x_{i_a^\star}(0), x_{i_b^\star}(0))$.

With Proposition 2, we know that the functions r_i with $i \in I_a^\star \sqcup I_b^\star$ are T-periodic, and variables x_i with $i \notin I_a^\star \sqcup I_b^\star$ converge asymptotically to 0. We thus consider the limiting dynamics

$$\begin{aligned} \dot{y}_a &= (\tilde{f}_a^l(t, y_a, y_b) - u(t))y_a, \\ \dot{y}_b &= (\tilde{f}_b^l(t, y_a, y_b) - u(t))y_b, \end{aligned} \tag{15}$$

where

$$\tilde{f}_a^l(t, y_a, y_b) := \tilde{\mu}_{i_a^\star}(s^l(t, y_a, y_b)), \quad \tilde{f}_b^l(t, y_a, y_b) := \tilde{\mu}_{i_b^\star}(s^l(t, y_a, y_b))$$

with

$$s^l(t, y_a, y_b) := s_{in} - \sum_{i \in I_a^\star} r_i(t)^{\alpha_i} y_a^{\frac{\alpha_i}{\alpha_{i^\star}}} - \sum_{i \in I_b^\star} r_i(t)^{\alpha_i} y_b^{\frac{\alpha_i}{\alpha_{i^\star}}},$$

which are time periodic functions. One has also

$$\lim_{t \to +\infty} |(\tilde{f}_a(t, y_a, y_b) - \tilde{f}_a^l(t, y_a, y_b))y_a| = \lim_{t \to +\infty} |(\tilde{f}_b(t, y_a, y_b) - \tilde{f}_b^l(t, y_a, y_b))y_b| = 0$$

uniformly for $(y_a, y_b) \in \mathcal{S}$. Therefore, the non-autonomous semi-flow of (14) in $\mathcal{S}$ is asymptotically periodic with limit periodic semi-flow of (15) (see Proposition 3.2 in [19]). The system (15) is competitive and we can apply the results of the literature about periodic competitive planar systems, which states that any bounded solution converges to a periodic solution $(y_a^p(\cdot), y_b^p(\cdot))$ (see for instance Theorem 4.2 in [15]). Let P be the Poincaré map associated to this dynamics

$$P : Y_0 \in \mathcal{S} \mapsto Y(T, Y_0) \in \mathcal{S}$$

where $Y(\cdot, Y_0)$ denotes the solution $(y_a(\cdot), y_b(\cdot))$ of (15) with $(y_a(0), y_b(0)) = Y_0$. On the axis $y_a = 0$ or $y_b = 0$, the dynamics is with a single taxon. One can then reproduce the arguments of the proof of Proposition 3 to show that there are unique fixed points $Y_a^\star = (y_a^\star, 0)$, $Y_b^\star = (0, y_b^\star)$ of P in $\mathcal{S}$ with $y_a^\star > 0$, $y_b^\star > 0$. Moreover, one has $y_a^\star = (x_a^p)_{i_a^\star}(0)$, $y_b^\star = (x_b^p)_{i_b^\star}(0)$, where $x_a^p(\cdot)$, $x_b^p(\cdot)$ are the asymptotic periodic solutions of (3) for the initial conditions ξ, ζ

$$\xi_i = \begin{cases} x_i(0), & i \in I_a, \\ 0, & i \in I_b, \end{cases} \qquad \zeta_i = \begin{cases} 0, & i \in I_a, \\ x_i(0), & i \in I_b \end{cases}$$

(remind from Proposition 2 that functions $\rho_{ii_a^\star}$ $(i \in I_a)$ or $\rho_{ii_b^\star}$ $(i \in I_b)$ remain the same).

The linearized dynamics $\dot{Y} = M(t)Y$ of (15) is given by the matrix

$$M(t) = \begin{bmatrix} a(t) & b(t) \\ c(t) & d(t) \end{bmatrix}$$

with

$$\begin{aligned} a(t) &= \tilde{f}_a^l(t, y_a(t), y_b(t)) - u(t) + \partial_{y_a} \tilde{f}_a^l(yt, y_a(t), y_b(t)) y_a(t), \\ b(t) &= \partial_{y_b} \tilde{f}_a^l(t, y_a(t), y_b(t)) y_a(t), \\ c(t) &= \partial_{y_a} \tilde{f}_b^l(t, y_a(t), y_b(t)) y_b(t), \\ d(t) &= \tilde{f}_b^l(t, y_a(t), y_b(t)) - u(t) + \partial_{y_b} \tilde{f}_b^l(t, y_a(t), y_b(t)) y_b(t)). \end{aligned}$$

Along the periodic solution $(x_a^p(\cdot), 0)$, one has

$$M(t) = \begin{bmatrix} \star & \star \\ 0 & f_b^l(t, ((x_a^p)_{i_a^\star}(t), 0) - u(t) \end{bmatrix}$$

for which the characteristics multiplier $\exp \int_0^T f_b^l(t, ((x_a^p)_{i_a^\star}(t), 0) - u(t)\, dt$ is equal to $\exp(T\lambda_{ba})$ and larger than 1 under condition (11). The fixed point $Y_a^\star$ is thus hyperbolic repulsive. In a similar way, $Y_b^\star$ is an hyperbolic repulsive fixed point under condition (12). This implies that $Y_a^\star$ and $Y_a^\star$ are isolated fixed points of P.

For $Y_0 = 0$, the solution of (15) is identically null and one has

$$M(t) = \begin{bmatrix} \mu_{i_a^\star}(s_{in}) - u(t) & 0 \\ 0 & \mu_{i_b^\star}(s_{in}) - u(t) \end{bmatrix}.$$

The characteristics multipliers are thus $\exp \int_0^T \mu_{i_a^\star}(s_{in}) - y(t)\, dt$, $\exp \int_0^T \mu_{i_b^\star}(s_{in}) - y(t)\, dt$ which are larger than one under Assumption 3. The zero solution is thus repulsive.

Along any positive solution $y_a(\cdot)$, $y_b(\cdot)$, note that one has $b(t) < 0$ and $c(t) < 0$ at any $t \in [0,T]$. Then, one has $\dot{Y}_1 > 0$ for $Y_1 = 0$ and $Y_2 < 0$, and $\dot{Y}_2 < 0$ for $Y_1 > 0$ and $Y_2 = 0$. Therefore, the second and fourth quadrant are invariant by the linear dynamics $\dot{Y} = M(t)Y$, which implies that the matrix $P'(Y_0)$ has strictly positive diagonal elements and strictly negative off-diagonal elements for a positive $Y_0 \in \mathcal{S}$. Let $\mathcal{S}' = \{Y \in \mathcal{S};\ s^l(0,Y) \in (\underline{s}, s_{in}]\}$. By Lemma 1, $\mathcal{S}'$ is invariant by P and any fixed point of P belongs to $\mathcal{S}'$, including $Y_a^\star$ and $Y_b^\star$. Following the arguments given in [2], the positive fixed points of P lie on a continuous curve Γ in $\mathcal{S}'$, which connects the fixed points $Y_a^\star$, $Y_b^\star$. Under Assumption 4, $x(\cdot)$ is analytic and the functions $r_i(\cdot)$ as well. Therefore, the map P is analytic on $\mathcal{S}'$. Then, the curve Γ is also analytic (see [2, 4]). If there were an infinite number of fixed points of P in $\mathcal{S}$ then all the points of the curve Γ will be fixed points by analyticity, which contradicts the fact that $Y_a^\star$ and $Y_b^\star$ are isolated fixed points.

Finally, as P has a finite number of fixed points on $\mathcal{S}$, we can apply the results about asymptotically autonomous discrete dynamical systems (Theorem 2.4 in [19]), from which one gets

$$\lim_{k\to+\infty} (x_{i_a^\star}(t+kT), x_{i_b^\star}(t+kT)) = Y(T, Y^\star), \quad t \in [0,T].$$

where $Y^\star$ is a fixed point of P in $\mathcal{S}$. As the fixed points on the axes 0, $Y_a^\star$ and $Y_b^\star$ are all repulsive, we conclude that $Y^\star$ is positive, and that $x(\cdot)$ converges asymptotically to the periodic solution $x^p(\cdot)$ given by

$$x_i^p(t) = \begin{cases} r_i(t)^{\alpha_i} y_a(t, Y^\star)^{\frac{\alpha_i}{\alpha_{i_a^\star}}} > 0, & i \in I_a^\star, \\ r_i(t)^{\alpha_i} y_b(t, Y^\star)^{\frac{\alpha_i}{\alpha_{i_b^\star}}} > 0, & i \in I_b^\star, \\ 0, & i \notin I_a^\star \sqcup I_b^\star, \end{cases}$$

for $t \in [0,T]$.

When $I_a^\star$ is not reduced to a singleton, take $i_a^\dagger \neq i_a^\star$ in $I_a^\star$ and consider perturbations $x^\varepsilon(\cdot)$ of the solution $x(\cdot)$, as solutions of (3) for the initial condition

$$x_i^\epsilon(0) = \begin{cases} x_i(0), & i \neq i_a^\dagger, \\ (r_{i_a^\dagger}(0) + \varepsilon)^{a_{i_a^\dagger}} x_{i_a^\star}(0)^{\frac{\alpha_i}{\alpha_{i_a^\star}}}, & i = i_a^\dagger, \end{cases} \tag{16}$$

with $\varepsilon > 0$. By continuity of solutions of (1) with respect to the initial condition, there exists $\bar{\varepsilon} > 0$ such that for any $\varepsilon \in (0, \bar{\varepsilon})$ $x^\epsilon(0)$ belongs to Δ and conditions (11), (12) are fulfilled for this new initial condition. As before, we deduce that $x^\epsilon(\cdot)$ converges asymptotically to a periodic solution $x^{\varepsilon,p}(\cdot)$ for which all species in $I_a^\star$ and $I_b^\star$ are present.

Let ρ_{ij}^ε be the ratio functions for the initial condition $x_i^\epsilon(0)$. Note from (16) that one has $\rho_{i_a^\dagger i_a^\star}^\varepsilon(0) = \rho_{i_a^\dagger i_a^\star}(0) + \varepsilon$. One gets from (7)

$$\rho_{i_a^\dagger i_a^\star}^\varepsilon(t) = \rho_{i_a^\dagger i_a^\star}^\varepsilon(0) \exp \int_0^t (\bar{u} - u(\tau)\big(\alpha_{i_a^\dagger}^{-1} - \alpha_{i_a^\star}^{-1}\big) d\tau = \rho_{i_a^\dagger i_a^\star}(t) + \varepsilon \exp \int_0^t (\bar{u} - u(\tau)\big(\alpha_{i_a^\dagger}^{-1} - \alpha_{i_a^\star}^{-1}\big) d\tau, \quad t \geq 0.$$

Therefore, the orbits $\gamma^+(\rho_{i_a^\dagger i_a^\star}^\varepsilon)$ for $\varepsilon \in (0, \bar{\varepsilon})$ are all distinct and we deduce, as in the proof of Proposition 3 that the orbits of the periodic solutions $\gamma^+(x^{\varepsilon,p})$ are all distinct.

□

Remark 2. *Conditions* (11), (12) *are independent of the choice of $i \in I_a^\star$, $I_b^\star$. Indeed, let $\bar{s}_a = \lambda_i(\bar{u})$ which is identical for any $i \in I_a^\star$, and one has*

$$\begin{array}{rcl} \langle \mu_i(s_b^p) \rangle_T - \bar{u} & = & \alpha_i \langle \varphi_a(s_b^p) \rangle_T + \beta_i - \bar{u} \\ & = & \alpha_i \langle \varphi_a(s_b^p) - \varphi_a(\bar{s}_a) \rangle_T + \alpha_i \varphi_a(\bar{s}_a) + \beta_i - \bar{u} \\ & = & \alpha_i \langle \varphi_a(s_b^p) - \varphi_a(\bar{s}_a) \rangle_T \end{array}$$

(using the property $\bar{u} = \mu_i(\bar{s}_a) = \alpha_i \varphi_a(\bar{s}_a) + \beta_i$ for $i \in I_a^\star$). The sign of $\langle \mu_i(s_b^p) \rangle_T - \bar{u}$ is thus independent of $i \in I_a^\star$, the numbers α_i being positive. One obtains symmetrically the same property for the sign of $\langle \mu_i(s_a^p) \rangle_T - \bar{u}$ with $i \in I_b^\star$.

Similarly to Proposition 3, Proposition 4 generalizes known results while giving indications about the number of periodic coexistence solutions of system (3). Indeed, when each taxon contains only one species, we obtain a result similar to the ones in [15]. However, thanks to our taxonomic assumptions, we are able to extend it from two species to two taxa, under very similar conditions. Then, we are also able to establish the existence of an infinite number of periodic coexistence solutions, generalizing the result of Proposition 3 from one to two taxa.

6 Numerical illustrations

We have considered three generating growth functions, among the class of Hill functions [8]

$$\varphi_a(s) = \frac{m_a s^2}{K_a^2 + s^2}, \qquad \varphi_b(s) = \frac{m_b s^4}{K_b^4 + s^4}, \qquad \varphi_c(s) = \frac{m_c s^2}{K_c^2 + s^2}. \tag{17}$$

with parameters given in Table 1, and operating conditions

$$s_{in} = 4, \quad \bar{u} = 0.8, \quad T = 10$$

where the periodic removal rate is

$$u(t) = \bar{u} + 0.2 \sin\left(\frac{2\pi t}{T}\right)$$

Graphs of the functions φ_a, φ_b, φ_c are depicted on Figure 1. Then, we have generated nine growth functions within these three taxa with characteristic numbers given in Table 2 and $\underline{s} = 0.5$. Let us denote the sets of indices of species belonging to a same taxon $I_a = \{1, 2, 3\}$, $I_b = \{4, 5, 6\}$, $I_c = \{7, 8, 9\}$. The graphs of these functions are depicted in Figure 2, where we have considered for each of these nine growth functions a C^1 extension for $s \leq \underline{s}$ as a polynomial increasing on $[0, \underline{s}]$ and null at 0. Clearly, Assumptions 1, 2, 3, 4 are satisfied. We have checked in all our simulations that the solutions remain in the set Δ for this value of $\underline{s}$. For the chosen value of $\bar{u}$, numbers $\bar{s}$ and γ_i defined in (5) in (6) are given in Table 3 for each taxon. Accordingly to Remark 1, the fittest species within each taxon are given by the subsets of indices $I_a^\star = \{1, 2\}$, $I_b^\star = \{4, 5\}$, $I_c^\star = \{8, 9\}$.

i	a	b	c
m_i	2	2.95	1.8
K_i	$\sqrt{3}$	$\sqrt{3}$	$\sqrt{3}$

Table 1: Parameters defining the generating functions φ_a, φ_b, φ_c

species	1	2	3	4	5	6	7	8	9
taxon	a			b			c		
α_i	0.9	1.15	0.85	0.8	1.05	0.6	0.9	1.1	0.7
β_i	0.161	−0.0165	0.1455	0.232	0.0545	0.326	0.161	0.019	0.261

Table 2: Characteristics numbers of the nine growth functions with respect to their taxon

species	1	2	3	4	5	6	7	8	9
taxon	a			b			c		
$\bar{s}$	1.28			1.30			1.40		
γ_i	0.09	0.09	0.03	0.09	0.09	0.01	0.09	0.09	0.03

Table 3: For each taxon, numbers $\bar{s}$ and γ_i (with $\bar{u} = 0.8$)

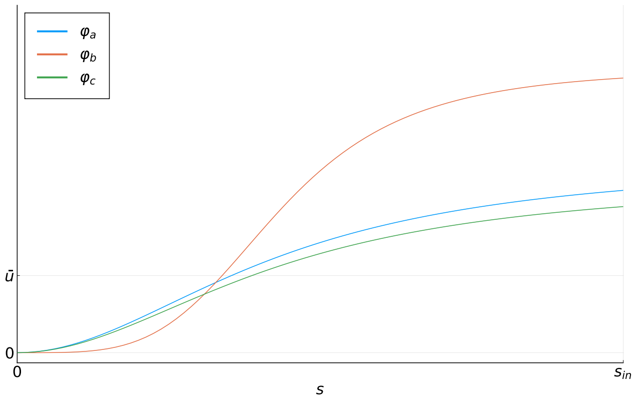
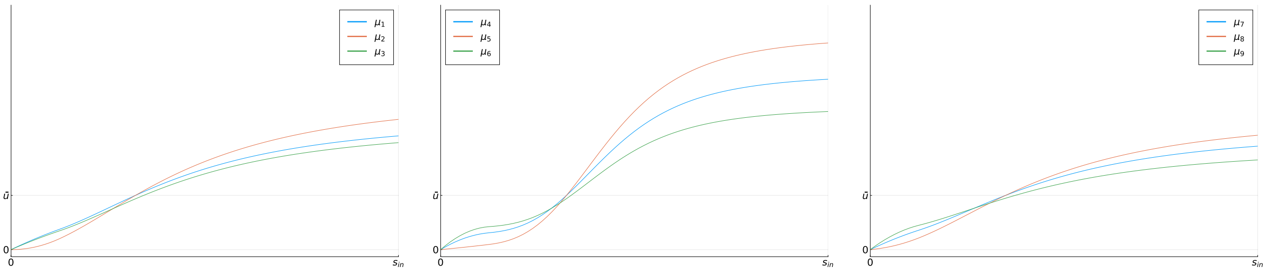


Figure 1: Graphs of the generating growth functions.

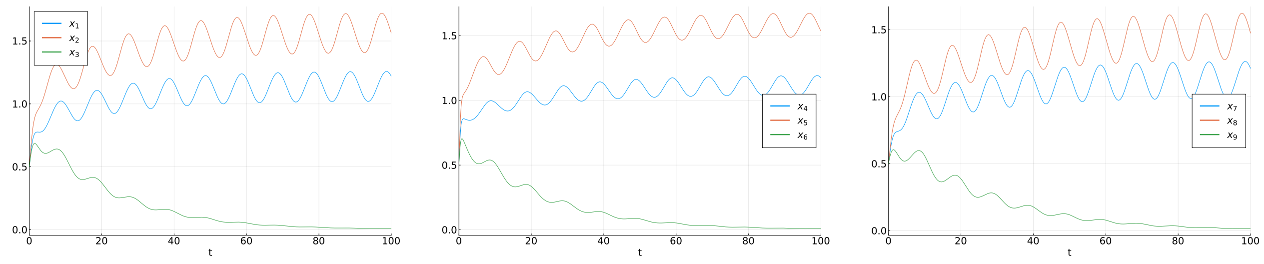


(a) Taxon I_a (b) Taxon I_b (c) Taxon I_c

Figure 2: Graphs of the growth functions for each taxon.

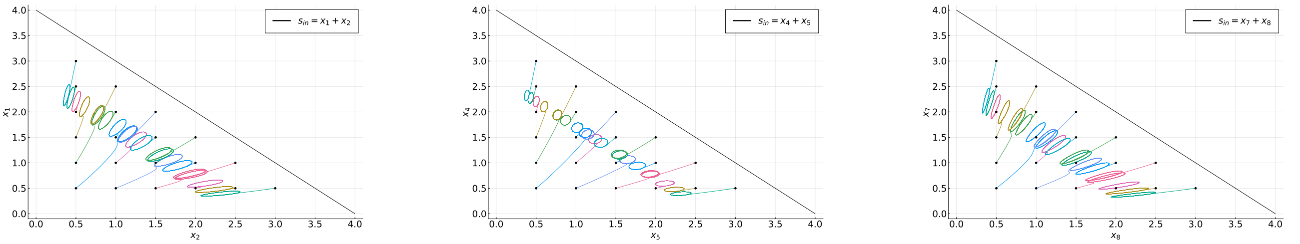
6.1 Simulations with species of a single taxon

One can first observe on Figure 3 that species that are not the fittest within their taxon i.e. that do not belong to $I_a^\star$, $I_b^\star$ or $I_c^\star$ are washed-out in presence of all the species in I_a, I_b or I_c respectively. This observation is in accordance with Proposition 3. Most interestingly, Figure 4 illustrates the multiplicity of periodic orbits with coexistence of two species, which is the main result of the present work.



(a) $x_i(\cdot)$ with $x_i(0) = 0.5$, for $i \in I_a$ (b) $x_i(\cdot)$ with $x_i(0) = 0.5$, for $i \in I_b$ (c) $x_i(\cdot)$ with $x_i(0) = 0.5$, for $i \in I_c$

Figure 3: Simulations with species of the same taxon.



(a) Orbits in the (x_1, x_2) plane when species in $I_a^\star$ only are present

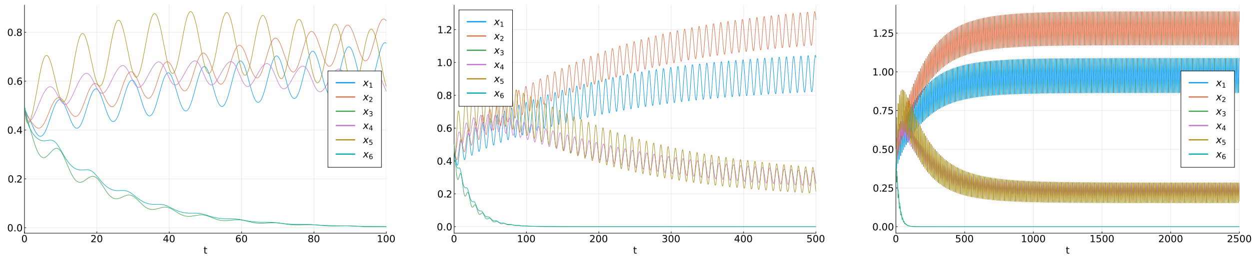
(b) Orbits in the (x_4, x_5) plane when species in $I_b^\star$ only are present

(c) Orbits in the (x_7, x_8) plane when species in $I_c^\star$ only are present

Figure 4: Multiplicity of periodic orbits among a single taxon.

6.2 Simulations with species of two different taxa

We have first considered species of I_a in presence with those of I_b (Figure 5). One can verify that these species satisfy the coexistence conditions given in Proposition 4 for the initial condition $(0.5, 0.5, 0.5, 0.5, 0.5, 0.5)$. Indeed, we computed numerically $\lambda_{ba} = 0.012576 > 0$ and $\lambda_{ab} = 0.006272 > 0$.



(a) run on the time interval $[0, 100]$

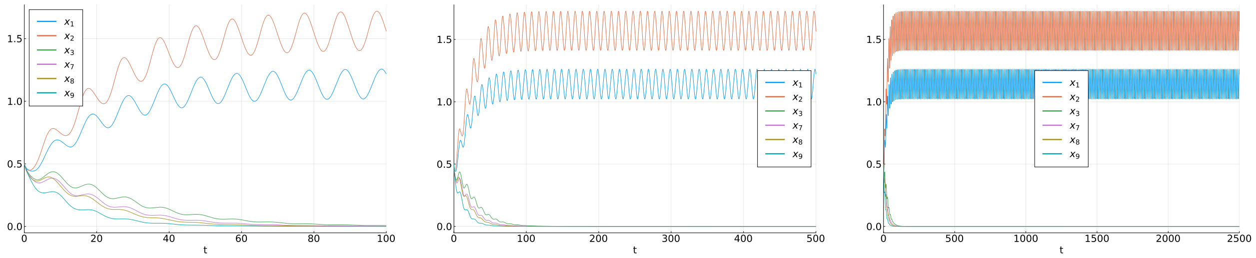
(b) run on the time interval $[0, 500]$

(c) run on the time interval $[0, 2500]$

Figure 5: Simulations of $x_i(\cdot)$ with $x_i(0) = 0.5$, for $i \in I_a \cup I_b$.

We observe that the species from $I_a \setminus I_a^\star$ and $I_b \setminus I_b^\star$ are washed-out while all the other species coexist. We also observe a transient behavior where the concentrations of species from $I_b^\star$ are initially raising faster than those of $I_a^\star$, before decreasing.

We have then considered species of I_a in presence with those of I_c (Figure 6). In this case, we computed $\lambda_{ca} = 0.071 > 0$ and $\lambda_{ac} = -0.0639 < 0$, which no longer guarantee the coexistence of the two taxa.



(a) run on the time interval $[0, 100]$

(b) run on the time interval $[0, 500]$

(c) run on the time interval $[0, 2500]$

Figure 6: Simulations of $x_i(\cdot)$ with $x_i(0) = 0.5$, for $i \in I_a \cup I_c$.

This time, all the species that do not belong to $I_a^\star$ are washed-out, as the second taxon is not well fit to survive in the competition.

6.3 Simulations under approximate taxon condition

The conditions for species to belong to a same taxon, and that some of them have identical numbers γ_i, are not generic among all increasing growth functions. In the spirit of former works [12, 3] for constant removal rate, we investigate here numerically cases where the condition (4) is only approximately satisfied.

For this purpose, we considered an additional species labeled $2'$ whose growth function $\mu_{2'}$ is closed to μ_2 but that does not belong to the taxon I_a. For the illustration, we have simply taken $\mu_{2'} = \mu_2 + \varepsilon\eta$, where η is a smooth function null at 0 with $\eta(\lambda_2(\bar{u})) > 0$ that is not proportionate to φ_a, and ε is a small number.

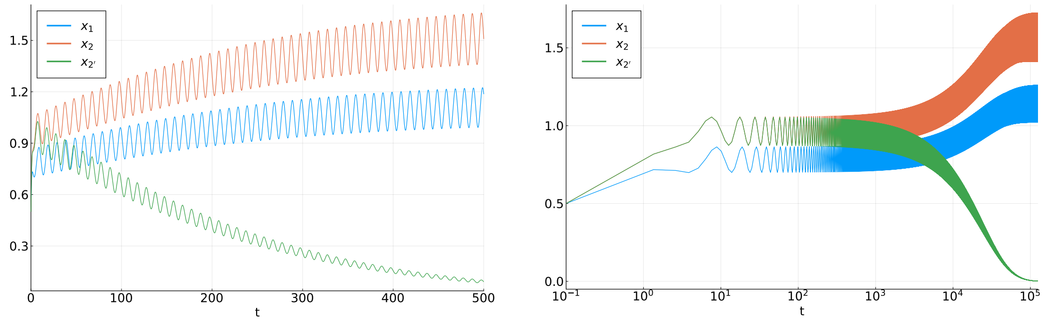
Remark 3. *In the particular case where the perturbation is such that $\mu_{2'}$ still satisfies the taxonomic assumption with $\alpha_{2'} = \alpha_2$ and $\gamma_{2'} = \gamma_2 - \epsilon$ (recall that $\gamma_2 = \gamma_1$), it is possible to study how the perturbation propagates over time. Indeed, we have from Proposition 1 that*

$$\dot{\rho}_{12'} = \left(\epsilon + (\bar{u} - u(t))(\alpha_1^{-1} - \alpha_2^{-1})\right) \rho_{12'}.$$

If $x_2(0) = x_{2'}(0)$ and the other initial conditions are the same, we can integrate this dynamics over $k \in \mathbb{N}$ periods from $t = 0$ and obtain the following comparison result:

$$\rho_{12'}(kT) = \rho_{12}(kT)e^{k\epsilon T}.$$

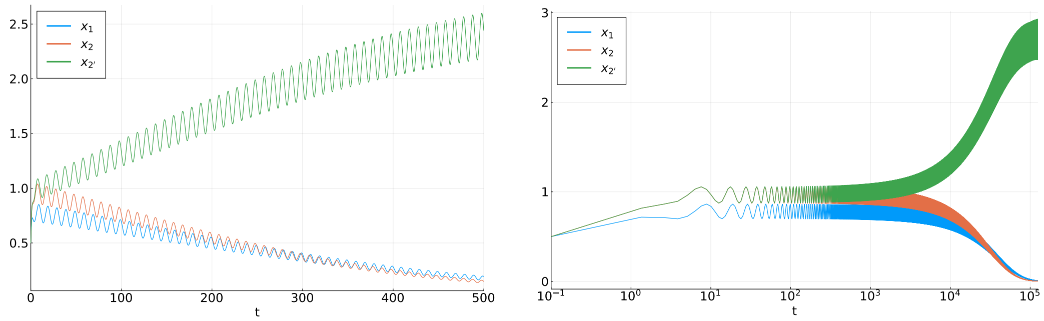
For the numerical illustration, we have taken $\eta(s) = \frac{s}{1+s}$ and first run simulations with species 1, 2 and $2'$. For $\varepsilon < 0$, the species $2'$ is asymptotically conducted to wash-out. However, as one can see on Figure 7, the transient can be very long when $\mu_{2'}$ is very close to μ_2 so that the three species coexist in an almost periodic manner during a long time horizon. On the opposite, for $\varepsilon > 0$, the species $2'$ is the final winner of the competition (Figure 8). However, it can take a long time for the other species to decline, so that here also the three species coexist in an almost periodic manner during a long time period.



(a) $\varepsilon = -10^{-2}$

(b) $\varepsilon = -10^{-4}$

Figure 7: Simulations with species 1, 2, $2'$ together when $\varepsilon < 0$.

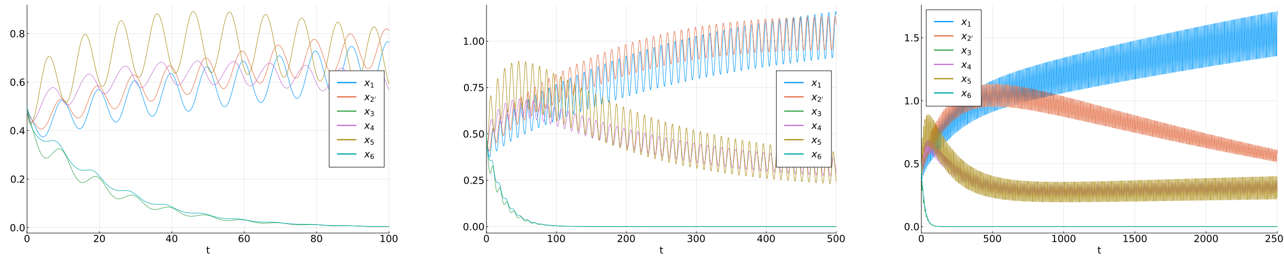


(a) $\varepsilon = 10^{-2}$

(b) $\varepsilon = 10^{-4}$

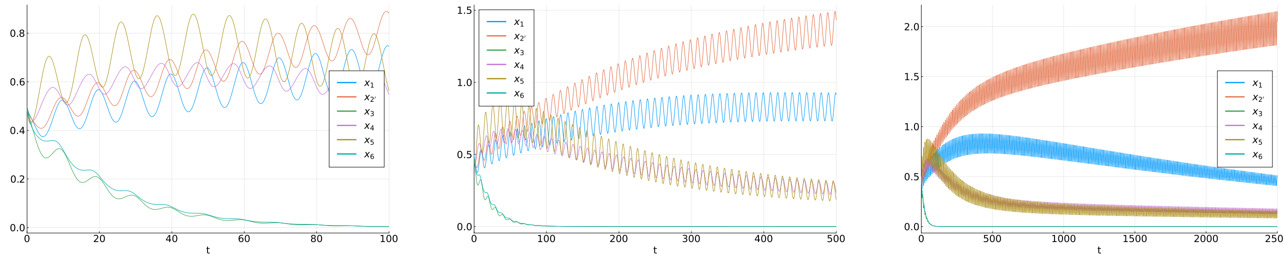
Figure 8: Simulations with species 1, 2, $2'$ together when $\varepsilon > 0$.

Then, we have considered species of taxa I_a and I_b together, as before (Figure 5), but where species 2 is replaced by species $2'$. Figures 9, 10 show that depending on the sign of ε, species $2'$ can be part or not of the final composition of the ecosystem, and that the time necessary to distinguish this issue can be very long, even for values of ε not extremely small. Note that coexistence of several species is maintained in both cases. A message here is that the analysis of non generic situations, as we do here, could be of some interest when facing cases likely to be closed from the non generic case, and we believe that this could be even more likely when considering many species.



(a) run on the time interval $t \in [0, 100]$ (b) run on the time interval $t \in [0, 500]$ (c) run on the time interval $t \in [0, 2500]$

Figure 9: Simulations with species 1, $2'$, 3, 4, 5, 6 together when $\varepsilon = -10^{-3}$.



(a) run on the time interval $[0, 100]$ (b) run on the time interval $[0, 500]$ (c) run on the time interval $[0, 2500]$

Figure 10: Simulations with species 1, $2'$, 3, 4, 5, 6 together when $\varepsilon = 10^{-3}$.

7 Conclusion

The present work has been motivated by the open problem about the number of periodic orbits that it is possible to observe in the multi-species chemostat model with periodic removal rate. We introduced conditions on subsets of species, that we call "taxa", and show with the help of a Poincaré map that the dynamics admits an infinite number of distinct periodic orbits when species belonging to a same taxon possess identical break-even concentrations. Moreover, we gave conditions to have coexistence between species of two different taxa, leading to a "double" infinity of periodic orbits. These constructions may serve as a counterexample to guide future constructions of sufficient conditions for uniqueness of periodic orbits in the chemostat model. Finally, we showed in simulation that although non-generic, these conditions could be "almost" met in practice leading to many almost periodic solutions with coexistence on a long time window.

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